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# TANGENT BUNDLES AND GAUGE GROUPS

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## Abstract

The differentials  $T^k a$   $(k \ge 1)$  of a diffeomorphism a of a smooth manifold M induce in the fibers of the fiber bundles  $T^k M$ , i.e., in the corresponding tangent spaces, linear transformations, which embody the action of the gauge group  $\mathcal{G}_k$ . This action extends in a natural way to the osculating subbundles  $\operatorname{Osc}^{k-1} M \subset T^k M$ .

 ${\bf Key}$  words: diffeomorphism of a smooth manifold, fiber bundles, action of the gauge group.

#### Introduction

The differential group  $\mathcal{G}$  of a smooth manifold M induces in the tangent bundle  $T^kM$  an action of the group of k-jets of transformations. More specifically, if a is a diffeomorphism of the manifold M, then its k-th differential  $T^ka$  is a transformation of the level  $T^kM$ . Then the level  $T^kM$  may be regarded as a homogeneous space  $J^k/H_k$ , where  $J^k$  is the group of k-jets of transformations and  $H_k$  is the stabilizer of an element  $u_{(u)} \in T^kM$ . The gauge group  $\mathcal{G}_k$  is defined as a certain subgroup of the linear group  $GL(2^kn,\mathbb{R})$ , where  $n = \dim M$  which is isomorphic to the stabilizer  $H_k$ . The action of the group  $\mathcal{G}$  extends to the osculating subbundle  $\operatorname{Osc}^{k-1}M \subset T^kM$ .

The paper contains all the necessary definitions and founds all the previous considerations. Commented examples and groups of derived formulas are presented as exercises.

## 1. Tangent groups

**1.1.** Leibniz rule. We apply the tangent functor T to the Cartesian product of smooth manifolds  $M_1$  and  $M_2$ :

$$T(M_1 \times M_2) = (TM_1 \times M_2) \oplus (M_1 \times TM_2),$$

and for the smooth mapping from  $M_1 \times M_2$  to some smooth manifold M

$$\lambda: M_1 \times M_2 \longrightarrow M: (u, v) \mapsto w = u \cdot v,$$

we define the tangent mapping  $T\lambda$ . First, by fixing the points  $u \in M_1$  and  $v \in M_2$  we define two mappings  $\lambda_u$  and  $\lambda_v$ :

$$\lambda_u: M_2 \to M: v \mapsto u \cdot v, \qquad \lambda_v: M_1 \to M: u \mapsto u \cdot v.$$

**Theorem 1.** To the pair of vectors  $u_1 \in T_uM_1$  and  $v_1 \in T_vM_2$  the mapping  $T\lambda$  associates the vector  $w_1 \in T_wM$ , and we have

$$w = u \cdot v \quad \Rightarrow \quad w_1 = u_1 \cdot v + u \cdot v_1 \,, \tag{1}$$

where  $u_1 \cdot v = T\lambda_v(u_1)$  and  $u \cdot v_1 = T\lambda_u(v_1)$ . In short, one can apply to the "product"  $w = u \cdot v$  the Leibniz rule.

**Proof.** We specify that, by means of the tangent maps  $T\lambda_v$  and  $T\lambda_u$ , two vectors  $u_1 \in T_u M_1$  and  $v_1 \in T_v M_2$  are transported from the points  $u \in M_1$  and  $v \in M_2$  to the point  $w \in M$ , where their sum defines the vector  $w_1 \in T_w M$ . Locally, this is confirmed by the formula:

$$w^{
ho} = \lambda^{
ho}(u^i, v^{lpha}) \quad \Rightarrow \quad w_1^{
ho} = \frac{\partial \lambda^{
ho}}{\partial u^i} u_1^i + \frac{\partial \lambda^{
ho}}{\partial v^{lpha}} v_1^{lpha},$$

where  $u^i, v^{\alpha}, w^{\rho}$  are the coordinates of the points u, v, w and  $u^i_1, v^{\alpha}_1, w^{\rho}_1$  are the components of the vectors  $u_1, v_1, w_1$  on the neighborhoods  $U_1 \subset M_1, U_2 \subset M_2, U \subset M, i = 1, \ldots, \dim M_1, \alpha = 1, \ldots, \dim M_2, \rho = 1, \ldots, \dim M$ .

Using the Leibniz rule we derive a set of important formulas in coordinate free form.

**Exercise 1: action of Leibniz rule.** Show that the Leibniz rule can be applied to the "product" of several factors, e.g.,

$$(u \cdot v \cdot w)_1 = u_1 \cdot v \cdot w + u \cdot v_1 \cdot w + u \cdot v \cdot w_1.$$

Exercise 2: prolongation of Leibniz rule. Prove that for the second tangent mapping  $T^2\lambda$ , the following formulas hold true:

$$= u \cdot v, \quad w_1 = u_1 \cdot v + u \cdot v_1, w_2 = u_2 \cdot v + u \cdot v_2, w_{12} = u_{12} \cdot v + u_2 \cdot v_1 + u_1 \cdot v_2 + u \cdot v_{12}.$$
(2)

Exercise 3: functional equation.

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Question: how can one solve the equation  $(u \cdot v)_1 = u_1 \cdot v + u \cdot v_1$  with respect to  $u_1$  for given  $v_1$  and  $(u \cdot v)_1$ , or relative to  $v_1$  for given  $u_1$  and  $(u \cdot v)_1$ ? This reminds the method of integration by parts:

$$d(uv) = u \, dv + v \, du \quad \rightsquigarrow \quad uv = \int u \, dv + \int v \, du, \quad \text{whence}$$
  
either  $\int u \, dv = uv - \int v \, du, \quad \text{or} \quad \int v \, du = uv - \int u \, dv.$ 

**1.2.** Coordinate-free story. The rule (1) is easy to use while building tangent groups and further, while studying representations of groups. If we have previously denoted the "product" of elements by a dot, as in (1) and (2), then while denoting the product of group elements, the dot will be omitted.

To a Lie group G with composition rule  $\gamma : (a, b) \mapsto c = ab$ , we associate the tangent group TG, having the composition law  $T\gamma$ :

$$c = ab \quad \Rightarrow \quad c_1 = a_1b + ab_1. \tag{3}$$

The vectors  $a_1 \in T_a G$  and  $b_1 \in T_b G$  are transported by means of the right shift  $r_b \doteq \gamma_b$ and of the left shift  $l_a \doteq \gamma_a$ , more exactly, by means of the tangent mappings  $Tr_b$  and  $Tl_a$ , from the points a and b to the point c, where the sum  $a_1b+ab_1 = Tr_b(a_1)+Tl_a(b_1)$ determines the vector  $c_1 \in T_c G$ . This is the composition law on the tangent group TG. The unity of the group TG is the null vector from  $T_e G$ . The inversion for the elements of TG is defined by the rule:

$$a_1 \in T_a G \quad \rightsquigarrow \quad a_1^{-1} = -a^{-1}a_1 a^{-1} \in T_{a^{-1}} G.$$
 (4)

**Exercise 4: unity and inverse elements.** Using (3), confirm the assertion regarding the unity of the group TG and the inversion of the elements (4). The formula (4) is obtained by solving the equation  $a_1b + ab_1 = 0$  relative to  $b_1$  for  $b = a^{-1}$ .

**Exercise 5: matrix representation.** Prove that the formulas (3) can be represented in matrix form as

$$\begin{pmatrix} a & 0 \\ a_1 & a \end{pmatrix} \cdot \begin{pmatrix} b & 0 \\ b_1 & b \end{pmatrix} = \begin{pmatrix} a & 0 \\ 0 & a \end{pmatrix} \cdot \begin{pmatrix} e & 0 \\ a^{-1}a_1 + b_1b^{-1} & e \end{pmatrix} \cdot \begin{pmatrix} b & 0 \\ 0 & b \end{pmatrix}.$$
 (5)

The sum of vectors in  $T_{ab}G$  reduces to the sum in  $TG_e$ :

c

$$a_1b + ab_1 = T(l_a \circ r_b)(a^{-1}a_1 + b_1b^{-1}).$$

Explain the meaning of the equality  $a_1b + ab_1 = (a_1a^{-1})c + c(b^{-1}b_1).$ 

**Exercise 6: second tangent group.** Prove that in the second tangent group  $T^2G$  the product of elements is defined by the formulas

$$= ab, \quad c_1 = a_1b + ab_1, \\ c_2 = a_2b + ab_2, \\ c_{12} = a_{12}b + a_2b_1 + a_1b_2 + ab_{12}.$$
(6)

and the inversion is performed by the rule

$$(a, a_1, a_2, a_{12})^{-1} \doteq (a^{-1}, a_1^{-1}, a_2^{-1}, a_{12}^{-1}), \text{ where}$$

$$a_1^{-1} = -a^{-1}a_1a^{-1},$$

$$a_2^{-1} = -a^{-1}a_2a^{-1},$$

$$a_{12}^{-1} = -a^{-1}a_{12}a^{-1} + a^{-1}a_2a^{-1}a_1a^{-1} + a^{-1}a_1a^{-1}a_2a^{-1}.$$
(7)

**Exercise 7: classical formulas.** Reduce the formulas (6) and (7) to the well known formulas from Analysis:

$$(uv)' = u'v + uv', \quad (uv)'' = u''v + 2u'v' + uv'' \left(\frac{1}{u}\right)' = -\frac{u'}{u^2}, \quad \left(\frac{1}{u}\right)'' = \frac{-uu'' + 2(u')^2}{u^3}.$$

Exercise 8: matrix relations. Using (6) and (7) prove the matrix relations:

$$\begin{pmatrix} c & 0 & 0 & 0 \\ c_1 & c & 0 & 0 \\ c_2 & 0 & c & 0 \\ c_{12} & c_2 & c_1 & c \end{pmatrix} = \begin{pmatrix} a & 0 & 0 & 0 \\ a_1 & a & 0 & 0 \\ a_2 & 0 & a & 0 \\ a_{12} & a_2 & a_1 & a \end{pmatrix} \cdot \begin{pmatrix} b & 0 & 0 & 0 \\ b_1 & b & 0 & 0 \\ b_2 & 0 & b & 0 \\ b_{12} & b_2 & b_1 & b \end{pmatrix}$$
$$\begin{pmatrix} a & 0 & 0 & 0 \\ a_1 & a & 0 & 0 \\ a_{12} & a_2 & a_1 & a \end{pmatrix}^{-1} = \begin{pmatrix} a^{-1} & 0 & 0 & 0 \\ a_1^{-1} & a^{-1} & 0 & 0 \\ a_2^{-1} & 0 & a^{-1} & 0 \\ a_{12}^{-1} & a_2^{-1} & a_1^{-1} & a^{-1} \end{pmatrix}.$$

Which endomorphism is involved here?

**Exercise 9: logarithmic derivatives.** Following the example of (5), represent the product of elements of the group  $T^2G$  in the form:

$$-\begin{pmatrix} a & 0 & 0 & 0 \\ a_1 & a & 0 & 0 \\ 0 & 0 & a & 0 \\ 0 & 0 & a_1 & a \end{pmatrix} \cdot \begin{pmatrix} e & 0 & 0 & 0 \\ 0 & e & 0 & 0 \\ a^{-1}a_2 + b_2b^{-1} & 0 & e & 0 \\ (a^{-1}a_2 + b_2b^{-1})_1 & a^{-1}a_2 + b_2b^{-1} & 0 & e \end{pmatrix} \cdot \begin{pmatrix} b & 0 & 0 & 0 \\ b_1 & b & 0 & 0 \\ 0 & 0 & b & 0 \\ 0 & 0 & b_1 & b \end{pmatrix},$$

where  $(a^{-1}a_2 + b_2b^{-1})_1 = a^{-1}a_{12} - a^{-1}a_1a^{-1}a_2 + b_{12}b^{-1} - b_2b^{-1}b_1b^{-1}$ . In this generalization, both logarithmic derivatives  $(\ln(uv))'$  and  $(\ln(uv))''$ , are present.

At the unity  $e \in G$  we fix the tangent vector  $e_1 \in T_eG$ . This vector is displaced by left shifts  $l_a$  over the group G to produce the left-invariant vector field  $ae_1$  and by right shifts  $r_a$ , to produce the right-invariant vector field  $e_1a$ . If at the unit  $e \in G$  we provide a frame, i.e., a basis of the space  $T_eG$ , then, in this way, two frame fields are defined on G (one right-invariant and left-invariant). The transition from one frame to another, at the point  $a \in G$ , is defined by some matrix A(a), which is an element of the linear group  $GL = GL(\dim G, \mathbb{R})$ . By this way, we define a homomorphism of the group G into the linear group GL:

$$G \to GL: a \mapsto A(a).$$
 (8)

Exercise 10: right/left shifts and inner automorphisms. Show that an 1-parametric subgroup  $a_t$  of the group G defines in the group G three flows corresponding to right shifts, left shifts and inner automorphisms

$$r_{a_t} = \exp tX, \quad l_{a_t} = \exp tX, \quad A_{a_t} = l_{a_t} \circ r_{a_t}^{-1} = \exp t(X - X)$$

and, accordingly, the left-invariant operator X, the right-invariant operator  $\hat{X}$  and the adjoint representation operator  $Y = \hat{X} - X$ . Prove this, using the formulas

$$Xf = (f \circ r_{a_t})'_{t=0}, \quad \widetilde{X}f = (f \circ l_{a_t})'_{t=0}, \quad Yf = (f \circ A_{a_t})'_{t=0}$$

where f is an arbitrary smooth function on G, taking into consideration that left shifts commute with right shifts.

**1.3.** Elements of representation theory. We consider a differentiable manifold, which we shall call *representation space* for the group G, or, in the following, simply *space*. A smooth mapping

$$\lambda : M \times G \longrightarrow M : (u, a) \mapsto v = u \cdot a$$

defines an *action* of the group G on the space M, if all the mappings

$$\lambda_a : M \to M \ u \mapsto u \cdot a , \quad \forall a \in G ,$$

are transformations (diffeomorphisms) of the space M, and the mapping  $a \mapsto \lambda_a$  is a homomorphism of the group G into the group of transformations of the space M. The homomorphism  $a \mapsto \lambda_a$  is understood either in the sense of the equality  $\lambda_{ab} = \lambda_a \circ \lambda_b$ or in the sense of the equality  $\lambda_{ab} = \lambda_b \circ \lambda_a$ . In the first case we say that the action of the group G on the space M is *left-sided*, and in the second case it is *right-sided*. By writing  $v = a \cdot u$  we have in view a left action, while by  $v = u \cdot a - a$  right action:

$$v = a \cdot u \quad \rightsquigarrow \quad (ab) \cdot u = a \cdot (b \cdot u),$$
  
$$v = u \cdot a \quad \rightsquigarrow \quad u \cdot (ab) = (u \cdot a) \cdot b.$$

The next formulas correspond to the right-sided action of the group G on the space M.

The kernel of the homomorphism  $a \mapsto \lambda_a$  is called the *stabilizer subgroup* of the group G. In the case of an effective action, the non-effectiveness kernel is trivial; it consists of the unity  $e \in G$ , and the mapping  $a \mapsto \lambda_a$  is injective.

For a fixed point  $u \in M$ , the mapping

$$\lambda_u : G \to M : a \mapsto u \cdot a$$

defines in the space M the *orbit* of this point. The whole space M is fibered into orbits  $\lambda_u(G)$ . When  $\lambda_u(G) = M$ , i.e., when the space M is the only orbit of the group G, we say that the action of the Lie group G on the space M is *transitive*. In such a case, M is called *homogeneous group space*. If moreover, the dimensions of G and M are equal, then the action of G on M is *simply transitive* and such an action defines an *exact representation* of the group G.

Equation 11: action by right/left shifts and inner automorphisms. Show that the actions of the group G on itself, provided by left and right shifts, are simply transitive actions. Prove that the actions provided by inner automorphisms are non-transitive.

The tangent map of the mapping  $\lambda$ , i.e.,  $T\lambda$ , defines a representation of the tangent group TG on the first level TM,

$$T\lambda$$
 :  $TM \times TG \rightarrow TM$  :  $(u_1, a_1) \mapsto v_1 = u_1 \cdot a + u \cdot a_1$ .

Formula (1) looks similarly:

$$v = u \cdot a \quad \rightsquigarrow \quad v_1 = u_1 \cdot a + u \cdot a_1. \tag{9}$$

We remark two particular cases:

for  $a_1 = 0$  we define the action of the group G on the level TM,

$$a_1 = 0 \quad \Rightarrow \quad u_1 \quad \mapsto \quad v_1 = u_1 \cdot a$$

for  $u_1 = 0$  we define the action of the tangent group TG on the space M,

$$u_1 = 0 \quad \Rightarrow \quad e_1 = a^{-1}a_1 \quad \mapsto \quad v_1 = u \cdot a_1 = v \cdot a^{-1}a_1 = v \cdot e_1.$$

The formula

$$v_1 = v \cdot a^{-1} a_1 \tag{10}$$

is the fundamental formula of the theory of Lie group representations.

In fact, to the vector  $e_1 = a^{-1}a_1 \in T_eG$  (which is an element of the Lie algebra g) at the point  $v \in M$  we associate some vector  $v_1 \in T_vM$ , and since v is an arbitrary point, it defines a vector field on the space M. This vector field is called the operator of the group G, or simply group operator<sup>1</sup>. Depending on the choice of the vector  $e_1 \in T_eG$ , in the space M we have an infinite set of group operators, and all of them, as vector fields, are tangent to the corresponding orbits.

For  $v_1 = 0$  the equality (10) provides the equation  $v \cdot a^{-1}a_1 = 0$ , or  $v \cdot e_1 = 0$ , which determines in the space  $T_e G$  those directions  $e_1$ , along which the point  $v \in M$  remains fixed. We define on the group G a Pfaff system, and its integral surface (solution), which contains the point  $e \in G$ , is a subgroup  $H_v \subset G$  called the stationary subgroup or the stabilizer of the point v.

In coordinates  $(v^{\alpha})$  on the neighborhood  $U \subset M$  of the point  $v \in M$ , Eq. (10) is written as a system  $dv^{\alpha} = \xi_i^{\alpha} \omega^i$ , where  $\omega^i$  are the forms of the left-invariant coframe on the group G. There appears a matrix<sup>2</sup>  $\xi = (\xi_i^{\alpha})$ , which determines a system of forms  $\vartheta^{\alpha}$  on the group G, and in the space M, a system of basic operators  $X_i$ :

$$\vartheta^{\alpha} = \xi_i^{\alpha} \omega^i, \quad X_i = \xi_i^{\alpha} \frac{\partial}{\partial v^{\alpha}}.$$

The number of operators  $X_i$  is equal to the dimension of G, and the number of forms  $\vartheta^{\alpha}$  is the dimension of M. Operators  $X_i$  and forms  $\vartheta^{\alpha}$  are not necessarily linearly independent. The Pfaff system  $\xi_i^{\alpha} \omega^i = 0$  for a fixed point  $v \in M$  is completely integrable and defines the stabilizer  $H_v \subset G$ .

**Exercise 12: vision from the classical theory.** Show that the system  $\xi_i^{\alpha} \omega^i = 0$  is the coordinate form of the equation  $v \cdot a^{-1}a_1 = 0$ .

**1.4.** Adjoint representation. In the groups G and TG we define the action by left shifts:

$$l_{a}: b \mapsto c = ab,$$
  

$$Tl_{a_{1}}: b_{1} \mapsto c_{1} = (a_{1}a^{-1})c + ab_{1},$$
  

$$\boxed{c_{1} = (a_{1}a^{-1})c},$$
(11)

<sup>&</sup>lt;sup>1</sup>Since the time of S. Lie and frequently nowadays, group operators have been called *infinitesimal* transformations or fundamental vector fields of the group.

transformations or fundamental vector fields of the group. <sup>2</sup>The matrix  $\xi$  plays an essential role in the theory of Lie group representations (see, e.g., S. Lie Theorems).

right shifts:

$$r_{a}: b \mapsto c = b a,$$
  

$$Tr_{a_{1}}: b_{1} \mapsto c_{1} = b_{1} a + c (a^{-1}a_{1}),$$
  

$$\boxed{c_{1} = c (a^{-1}a_{1})},$$
(12)

and inner automorphisms:

$$A_{a}: b \mapsto c = aba^{-1},$$
  

$$TA_{a_{1}}: b_{1} \mapsto c_{1} = (a_{1}a^{-1})c - c(a_{1}a^{-1}) + ab_{1}a^{-1},$$
  

$$\boxed{c_{1} = (a_{1}a^{-1})c - c(a_{1}a^{-1})}.$$
(13)

The basic formula (10) is rewritten, for  $b_1 = 0$ , in the forms (11), (12) and (13), respectively.

Inner automorphisms are directly related to higher order movements.

Hence, when in the spaces A and B there take place the transformations a and b, the mapping  $\varphi: A \to B$  is brought into the mapping  $\tilde{\varphi}: A \to B$ . This is shown by the diagram:

If we set here A = B, a = b and if  $\varphi$  is a diffeomorphism, i.e.,  $\varphi$  is a transformation of the space A, then this diagram describes the transformation of the mapping  $\varphi$ , subject to the influence of the transformation a:

The transformation  $\varphi$  is subject to the inner automorphism.

**Exercise 13: higher order transformations.** The transformation of order 2  $\varphi \rightsquigarrow \tilde{\varphi}$  is described by the 2-dimensional diagram (1.14). Show that the transformation of order 3, i.e., a transformation of transformation  $\varphi \rightsquigarrow \tilde{\varphi}$ , is described by a 3-dimensional diagram and the transformation of order k is described by a corresponding k-dimensional diagram.

If the arrow a in diagram (14) is assumed to represent the 1-parametric group  $a_t$  of transformations of the space A, or in brief, the flow  $a_t$ , then we see how, to a change of the parameter t (of time), it corresponds to a change of the mapping  $\varphi_t = a_t \varphi a_t^{-1}$ . We can talk then about a 1-parametric family of mappings  $\varphi_t$  in the field  $a_t$ .

If the arrow  $\varphi$  in diagram (14) is regarded as a 1-parametric group of transformations  $b_{\tau}$  of the space A (the flow  $b_{\tau}$ ), then we can see how this flow changes under the transformation a, i.e.,  $b_{\tau} \rightsquigarrow \tilde{b}_{\tau} = ab_{\tau}a^{-1}$ .

**Exercise 14: transformation of the flow.** Show that if  $b_{\tau}$  is the flow of the vector field Y and  $\tilde{b}_{\tau}$  is the flow of the field  $\tilde{Y}$ , then  $\tilde{Y} = TaY$ , and the tangent mapping Ta acts on the field Y:

$$b_{\tau} = \exp \tau Y \quad \rightsquigarrow \quad \widetilde{b}_{\tau} = ab_{\tau}a^{-1} = \exp \tau \widetilde{Y}, \qquad Y \rightsquigarrow \widetilde{Y} = TaY.$$

**Exercise 15: interaction of vector fields.** Let X and Y be two vector fields. The flows of these fields  $a_t = \exp tX$  and  $b_\tau = \exp \tau Y$  interact according to the scheme:

$$\begin{aligned} b_{\tau} & \rightsquigarrow & a_t b_{\tau} a_{-t}, & Y \rightsquigarrow T a_t Y, \\ a_t & \rightsquigarrow & b_{\tau} a_t b_{-\tau}, & X \rightsquigarrow T b_{\tau} X. \end{aligned}$$

Using derivatives of the function f

$$Xf = (f \circ a_t)'_{t=0}$$
 and  $Yf = (f \circ b_\tau)'_{\tau=0}$ 

perform the differentiation (the parameters t or  $\tau$  from above the arrow mean differentiation relative to t for  $t \to 0$  or to  $\tau$ , for  $\tau \to 0$ ):

$$\begin{aligned} f \circ (a_t b_\tau a_t^{-1}) & \stackrel{t}{\longrightarrow} & (Xf) \circ b_\tau - X(f \circ b_\tau) & \stackrel{\tau}{\longrightarrow} & (YX - XY)f \\ f \circ (b_\tau a_t b_\tau^{-1}) & \stackrel{\tau}{\longrightarrow} & (Yf) \circ a_t - Y(f \circ a_t) & \stackrel{t}{\longrightarrow} & (XY - YX)f. \end{aligned}$$

Check the validity of the relation  $(Ta_tY)'_{t=0} = -(Tb_{\tau}X)'_{\tau=0}$  and establish a connection with the brackets [X, Y] = XY - YX.

If in one flow the points move along trajectories and under the influence of the other flow, this movement is transformed, and then the movement of the movement takes place, or a second-order movement. Under the influence of a third flow, the movement of second order changes its shape, then the movement of third order occurs, etc. In the infinitesimal approach this reduces to the iterations

$$a_t \rightsquigarrow Ta_t \rightsquigarrow T^2a_t \rightsquigarrow \ldots$$

and to the corresponding vector fields on the levels

$$X \rightsquigarrow \overset{(1)}{X} \rightsquigarrow \overset{(2)}{X} \rightsquigarrow \dots, \qquad T^k a_t = \exp t \overset{(k)}{X}, \quad k = 0, 1, 2, \dots$$
(15)

In this way, the flow  $T^k a_t$  induces a movement of order k.

1.5. Gauge groups. Let us fix on each level a point

$$u_{(k)} \in T^k M$$
, such that  $\pi_k(u_{(k)}) = u_{(k-1)}, \ k = 0, 1, 2, \dots$ 

In the neighborhood  $T^k U \subset T^k M$  these points are defined by their coordinates:

$$U : u_{(0)} \rightsquigarrow (u^{i}),$$
  

$$TU : u_{(1)} \rightsquigarrow (u^{i}, u^{i}_{1}),$$
  

$$T^{2}U : u_{(2)} \rightsquigarrow (u^{i}, u^{i}_{1}, u^{i}_{2}, u^{i}_{12}),$$
  

$$T^{3}U : u_{(3)} \rightsquigarrow (u^{i}, u^{i}_{1}, u^{i}_{2}, u^{i}_{12}, u^{i}_{3}, u^{i}_{13}, u^{i}_{23}, u^{i}_{123}),$$
  
..........

Transformation of coordinates in the neighborhood  $U \subset M$ 

$$u^i \rightsquigarrow \widetilde{u}^i \circ a = a^i$$

induces a change of coordinates in each neighborhood  $T^k U$ :

$$(u^{i}, u^{i}_{1}, u^{i}_{2}, u^{i}_{12}, \dots) \iff (\widetilde{u}^{i}, \widetilde{u}^{i}_{1}, \widetilde{u}^{i}_{2}, \widetilde{u}^{i}_{12}, \dots) = (a^{i}, a^{i}_{1}, a^{i}_{2}, a^{i}_{12}, \dots).$$

Namely, if these transformations of coordinates in the neighborhood TU are defined by the system

$$\begin{cases} \widetilde{u}^i = a^i, \\ \widetilde{u}^i_1 = a^i_1 \doteq a^i_j u^j_1, \end{cases}$$

with the Jacobian block-matrix

$$\begin{pmatrix} a_j^i & 0\\ (a_j^i)_1 & a_j^i \end{pmatrix}, \quad \text{where} \quad a_j^i = \frac{\partial a^i}{\partial u^j}, \quad a_{jk}^i = \frac{\partial^2 a^i}{\partial u^j \partial u^k}, \quad (a_j^i)_1 \doteq a_{jk}^i u_1^k,$$

then the transformation of coordinates in the neighborhood  $T^2U$  are defined by the system

$$\begin{cases} u^{i} = a^{i}, \\ \widetilde{u}_{1}^{i} = a_{1}^{i} \doteq a_{j}^{i}u_{1}^{j}, \\ \widetilde{u}_{2}^{i} = a_{2}^{i} \doteq a_{j}^{i}u_{2}^{j}, \\ \widetilde{u}_{12}^{i} = a_{12}^{i} \doteq a_{jk}^{i}u_{1}^{j}u_{2}^{k} + a_{j}^{i}u_{12}^{j}, \end{cases}$$

with the Jacobian block-matrix

$$\begin{pmatrix} a_j^i & 0 & 0 & 0 \\ (a_j^i)_1 & a_j^i & 0 & 0 \\ (a_j^i)_2 & 0 & a_j^i & 0 \\ (a_j^i)_{12} & (a_j^i)_2 & (a_j^i)_1 & a_j^i \end{pmatrix}, \quad \text{where} \quad \begin{cases} (a_j^i)_1 \doteq a_{jk}^i u_1^k, \\ (a_j^i)_2 \doteq a_{jk}^i u_2^k, \\ (a_j^i)_{12} \doteq a_{jkl}^i u_1^k u_2^l + a_{jk}^i u_{12}^k, \end{cases}$$

 $\operatorname{etc.}$ 

When performing a lift from one level to another,  $U \rightsquigarrow TU \rightsquigarrow T^2U \rightsquigarrow \ldots$ , the Jacobian matrix is inductively built according to the scheme:

$$\mathfrak{a} \rightsquigarrow \begin{pmatrix} \mathfrak{a} & 0\\ \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix} \rightsquigarrow \begin{pmatrix} \mathfrak{a} & 0 & 0 & 0\\ \mathfrak{a}_1 & \mathfrak{a} & 0 & 0\\ \mathfrak{a}_2 & 0 & \mathfrak{a} & 0\\ \mathfrak{a}_{12} & \mathfrak{a}_2 & \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix} \rightsquigarrow \dots,$$
(16)

with repeated, as shown above, n-dimensional blocks

$$\mathfrak{a} = (a_j^i), \ \mathfrak{a}_1 = (a_j^i)_1, \ \mathfrak{a}_2 = (a_j^i)_2, \ \mathfrak{a}_{12} = (a_j^i)_{12}, \ldots$$

Therefore, there follows the general rule: the Jacobian matrix of the transformation of coordinates on the neighborhood  $T^kU$  is of the form

$$\begin{pmatrix} \mathcal{A} & 0\\ \mathcal{A}_k & \mathcal{A} \end{pmatrix},\tag{17}$$

where the block  $\mathcal{A}$  is the Jacobian matrix on  $T^{k-1}U$  and  $\mathcal{A}_k \doteq d_k \mathcal{A}, k = 1, 2, ...$ 

In other words, the Jacobian matrix on the neighborhood  $T^k U$  consists of four blocks, where the Jacobian matrix  $\mathcal{A}$  of the neighborhood  $T^{k-1}U$  is repeated on the diagonal, the upper-right block is zero, and the left-lower block is the differential of the block  $\mathcal{A}$  taking into consideration the k-th level, i.e.,  $\mathcal{A}_k \doteq d_k \mathcal{A}$ .

Formula (17) defines the sequence of matrices (16).

**Exercise 16: inversion rule.** Show that the inversion of the matrix (17) takes place according to the scheme: (4, -2) = (-2, -2) = (4, -2)

$$\begin{pmatrix} \mathcal{A} & 0 \\ \mathcal{A}_{k} & \mathcal{A} \end{pmatrix} = \begin{pmatrix} \mathcal{E} & 0 \\ \mathcal{A}_{k}\mathcal{A}^{-1} & \mathcal{E} \end{pmatrix} \cdot \begin{pmatrix} \mathcal{A} & 0 \\ 0 & \mathcal{A} \end{pmatrix} \rightsquigarrow$$
$$\rightarrow \quad \begin{pmatrix} \mathcal{A} & 0 \\ \mathcal{A}_{k} & \mathcal{A} \end{pmatrix}^{-1} = \begin{pmatrix} \mathcal{A}^{-1} & 0 \\ 0 & \mathcal{A}^{-1} \end{pmatrix} \cdot \begin{pmatrix} \mathcal{E} & 0 \\ -\mathcal{A}_{k}\mathcal{A}^{-1} & \mathcal{E} \end{pmatrix},$$

where  $\mathcal{E}$  is the identity block. See (4) and

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$$a_1 = (a_1 a^{-1})a \iff a_1^{-1} = -a^{-1}(a_1 a^{-1}).$$

The matrix (17) depends on the point  $u_{(k)} \in T^k U$ . If this point is fixed, then a numeric matrix is defined, but still having the freedom to choose the function  $a^i$ (or the corresponding jet of the transformation). **Exercise 17: gauge group.** Show that all matrices of the form (17), with the point  $u_{(k)} \in T^k U$  fixed, determine a subgroup of the linear group of order  $2^k n$ ,

$$\mathcal{G}_k \subset GL(2^k n, \mathbb{R}).$$

Prove the existence of the groups  $\mathcal{G}_1$ ,  $\mathcal{G}_2$ ,  $\mathcal{G}_3$  and extend to  $\mathcal{G}_k$ .

We call the group  $\mathcal{G}_k$  of matrices (17) with fixed point  $u_{(k)} \in T^k U$  the gauge group of order k on the manifold M. By setting  $k = 0, 1, 2, 3, \ldots$ , we obtain an infinite sequence of gauge groups

$$\mathcal{G} \rightsquigarrow \mathcal{G}_1 \rightsquigarrow \mathcal{G}_2 \rightsquigarrow \mathcal{G}_3 \rightsquigarrow \dots$$
 (18)

**Theorem 2.** The gauge group of order k is isomorphic to the k-th tangent group of the linear group  $GL(n,\mathbb{R})$ , which, in its turn, is embedded in the linear group  $GL(2^kn,\mathbb{R})$ :

$$\mathcal{G}_k \approx T^k \big( GL(n, \mathbb{R}) \big) \subset GL(2^k n, \mathbb{R}), \quad k = 0, 1, 2, \dots$$
(19)

In this case

$$\dim GL(2^k n, \mathbb{R}) = (2^k n)^2 \quad and \quad \dim \mathcal{G}_k = 2^k n^2.$$

**Proof.** We fix the element  $u_{(k)} \in T^k M$  of the k-th level. Matrices (17) generate a subgroup  $\mathcal{G}_k$  of the linear group  $GL(2^k n, \mathbb{R})$  (see Exercise 17). The fixing of the point  $u_{(k)}$  does not limit the freedom of choice for the element (17) in the group  $\mathcal{G}_k$ . Hence, the group  $\mathcal{G}_k$  is uniquely defined regardless of the point  $u_{(k)} \in T^k M$ . On the other side, the tangent group  $T^k(GL(n,\mathbb{R}))$  coincides up to an isomorphism, with the matrix group (17), or  $\mathcal{G}_k$ . This follows from the formulas (3)–(7) and Exercises 12–17, if we assume  $G = GL(n,\mathbb{R})$ .

Further, in the matrix (17), besides the point  $u_{(k)} \in T^k U$ , there exists the k-jet of coordinate transformations  $(a_i^j, a_{i_1i_2}^j, \ldots, a_{i_1i_2\dots i_k}^j)$ . We shall denote as  $J_k$  the group of such jets at the point u (see Exercise 5). The homomorphism is defined:

$$\chi_k: J_k \to \mathcal{G}_k. \tag{20}$$

**Exercise 18: jets and gauge group.** Show that for k = 2, the mapping  $\chi_2$  is homomorphic, i.e., to a composition of 2-jets  $(a_k^i, a_{kl}^j)$  and  $(b_l^i, a_{kl}^i)$  there corresponds the product of matrices  $\mathcal{A}_2$ ,

$$\begin{pmatrix} a_k^i & 0 \\ a_k^i u_1^l & a_k^i \end{pmatrix} \cdot \begin{pmatrix} b_j^k & 0 \\ b_j^k u_1^l & b_j^k \end{pmatrix} = \begin{pmatrix} a_k^i b_j^k & 0 \\ (a_k^i b_j^k)_l u_1^l & a_k^i b_j^k \end{pmatrix},$$

and to the inverse 2-jet  $(a_j^i, a_{jl}^i)^{-1} \doteq (\bar{a}_j^i, -\bar{a}_s^i a_{kl}^s \bar{a}_j^k)$ , there corresponds the inverse matrix  $\mathcal{A}_2^{-1}$ 

$$\begin{pmatrix} a_k^i & 0 \\ a_k^i u_1^l & a_k^i \end{pmatrix}^{-1} = \begin{pmatrix} \bar{a}_j^i & 0 \\ -\bar{a}_s^i a_{kl}^s \bar{a}_j^k u_1^l & \bar{a}_j^i \end{pmatrix}$$

Generalize this to the general case k.

Exercise 19: homogeneity of tangent space. Show that the kernel of the homomorphism  $\chi_k$  is the stabilizer  $H_{u_{(k)}}$  of the element  $u_{(k)} \in T^k M$  in the group  $J_k$ . The tangent space  $T^k_{u_{(k)}} M$  is identified with the homogeneous space  $J_k/H_{u_{(k)}}$ .

Let us consider once again the gauge groups of the sequence (18). The first group  $\mathcal{G}$  is the linear group  $GL(n, \mathbb{R})$ ,

$$\mathcal{G} = GL(n, \mathbb{R}).$$

The second group  $\mathcal{G}_1$  is isomorphic to the tangent group  $T(GL(n,\mathbb{R}))$ . Its elements are block matrices of the form

$$\begin{pmatrix} \mathfrak{a} & 0\\ \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix}$$
, where  $\mathfrak{a} \in GL(n, \mathbb{R})$  and  $\mathfrak{a}_1 \in gl(n, \mathbb{R})$ .

The correspondence  $\mathcal{G}_1 \iff T(GL(n,\mathbb{R}))$  is one-to-one. The product of elements in the group  $\mathcal{G}_1$ ,

$$\begin{pmatrix} \mathfrak{a} & 0 \\ \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix} \cdot \begin{pmatrix} \mathfrak{b} & 0 \\ \mathfrak{b}_1 & \mathfrak{b} \end{pmatrix} = \begin{pmatrix} \mathfrak{a} \mathfrak{b} & 0 \\ (\mathfrak{a} \mathfrak{b})_1 & \mathfrak{a} \mathfrak{b} \end{pmatrix},$$

reduces to the Leibniz rule in the tangent group  $T(GL(n,\mathbb{R}))$ ,

$$(\mathfrak{ab})_1 = \mathfrak{a}_1\mathfrak{b} + \mathfrak{ab}_1,$$

and the inversion of elements in  $\mathcal{G}_1$ ,

$$\begin{pmatrix} \mathfrak{a} & 0 \\ \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix}^{-1} = \begin{pmatrix} \mathfrak{a}^{-1} & 0 \\ -\mathfrak{a}^{-1}\mathfrak{a}_1\mathfrak{a}^{-1} & \mathfrak{a}^{-1} \end{pmatrix},$$

reduces to the rule

$$\mathfrak{a}_1^{-1} = -\mathfrak{a}^{-1}\mathfrak{a}_1\mathfrak{a}^{-1}.$$

This speaks about an isomorphism between the groups  $\mathcal{G}_1$  and  $T(GL(n, \mathbb{R}))$ . An inner authomorphism in  $\mathcal{G}_1$  is generated as allows:

$$\begin{pmatrix} \mathfrak{a} & 0 \\ \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix} \cdot \begin{pmatrix} \mathfrak{b} & 0 \\ \mathfrak{b}_1 & \mathfrak{b} \end{pmatrix} \cdot \begin{pmatrix} \mathfrak{a} & 0 \\ \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix}^{-1} = \begin{pmatrix} \mathfrak{a}\mathfrak{b}\mathfrak{a}^{-1} & 0 \\ (\mathfrak{a}\mathfrak{b}\mathfrak{a}^{-1})_1 & \mathfrak{a}\mathfrak{b}\mathfrak{a}^{-1} \end{pmatrix},$$

with the block  $(\mathfrak{a}\mathfrak{b}\mathfrak{a}^{-1})_1 = \mathfrak{a}\mathfrak{b}_1\mathfrak{a}^{-1} + \mathfrak{a}_1\mathfrak{a}^{-1}(\mathfrak{a}\mathfrak{b}\mathfrak{a}^{-1}) - (\mathfrak{a}\mathfrak{b}\mathfrak{a}^{-1})\mathfrak{a}_1\mathfrak{a}^{-1},$  etc.

The following group  $\mathcal{G}_2$  is isomorphic to the tangent group  $T^2(GL(n,\mathbb{R}))$ . The stair-like structure appears again:

$$\begin{pmatrix} \mathfrak{a} & 0 & 0 & 0 \\ \mathfrak{a}_1 & \mathfrak{a} & 0 & 0 \\ \mathfrak{a}_2 & 0 & \mathfrak{a} & 0 \\ \mathfrak{a}_{12} & \mathfrak{a}_2 & \mathfrak{a}_1 & \mathfrak{a} \end{pmatrix} \cdot \begin{pmatrix} \mathfrak{b} & 0 & 0 & 0 \\ \mathfrak{b}_1 & \mathfrak{b} & 0 & 0 \\ \mathfrak{b}_2 & 0 & \mathfrak{b} & 0 \\ \mathfrak{b}_{12} & \mathfrak{b}_2 & \mathfrak{b}_1 & \mathfrak{b} \end{pmatrix} = \begin{pmatrix} \mathfrak{a} \mathfrak{b} & 0 & 0 & 0 \\ (\mathfrak{a} \mathfrak{b})_1 & \mathfrak{a} \mathfrak{b} & 0 & 0 \\ (\mathfrak{a} \mathfrak{b})_2 & 0 & \mathfrak{a} \mathfrak{b} & 0 \\ (\mathfrak{a} \mathfrak{b})_{12} & (\mathfrak{a} \mathfrak{b})_2 & (\mathfrak{a} \mathfrak{b})_1 & \mathfrak{a} \mathfrak{b} \end{pmatrix},$$

where

$$\begin{split} (\mathfrak{a}\mathfrak{b})_1 &= \mathfrak{a}_1\mathfrak{b} + \mathfrak{a}\mathfrak{b}_1, \\ (\mathfrak{a}\mathfrak{b})_2 &= \mathfrak{a}_2\mathfrak{b} + \mathfrak{a}\mathfrak{b}_2, \\ (\mathfrak{a}\mathfrak{b})_{12} &= \mathfrak{a}_{12}\mathfrak{b} + \mathfrak{a}_2\mathfrak{b}_1 + \mathfrak{a}_1\mathfrak{b}_2 + \mathfrak{a}\mathfrak{b}_{12}. \end{split}$$

Exercise 20: logarithmic rule for gauge group. Show that while forming the blocks

$$\mathfrak{a} \rightsquigarrow \mathfrak{a}^{-1}\mathfrak{a}_1 \rightsquigarrow (\mathfrak{a}^{-1}\mathfrak{a}_1)_2 = \mathfrak{a}^{-1}\mathfrak{a}_{12} - \mathfrak{a}^{-1}\mathfrak{a}_2\mathfrak{a}^{-1}\mathfrak{a}_1 \rightsquigarrow \dots$$

there appears the following property of the logarithmic function:

$$\ln u \iff \frac{u'}{u} \iff \frac{u''}{u} - \frac{(u')^2}{u^2} \iff \dots$$

We shall further denote the Lie algebra of the group  $\mathcal{G}_k$  by  $\overline{\mathcal{G}}_k$ .

The general scheme is the following. An element of the group  $\mathcal{G}_k$  is generated according to the principle:

$$\begin{pmatrix} \mathcal{A} & 0 \\ \mathcal{A}_k & \mathcal{A} \end{pmatrix}, \quad \text{where} \quad \mathcal{A} \in \mathcal{G}_{k-1}, \ \mathcal{A}_k \in \overline{\mathcal{G}}_{k-1}.$$

The product and the inversion of elements,

$$\begin{pmatrix} \mathcal{A} & 0 \\ \mathcal{A}_k & \mathcal{A} \end{pmatrix} \cdot \begin{pmatrix} \mathcal{B} & 0 \\ \mathcal{B}_k & \mathcal{B} \end{pmatrix} = \begin{pmatrix} \mathcal{A}\mathcal{B} & 0 \\ (\mathcal{A}\mathcal{B})_k & \mathcal{A}\mathcal{B} \end{pmatrix},$$
$$\begin{pmatrix} \mathcal{A} & 0 \\ \mathcal{A}_k & \mathcal{A} \end{pmatrix}^{-1} = \begin{pmatrix} \mathcal{A}^{-1} & 0 \\ \mathcal{A}_k^{-1} & \mathcal{A}^{-1} \end{pmatrix},$$

reduce to the rules:

$$(\mathcal{AB})_k = \mathcal{A}_k \mathcal{B} + \mathcal{AB}_k, \quad \mathcal{A}_k^{-1} = -\mathcal{A}^{-1} \mathcal{A}_k \mathcal{A}^{-1}$$

The Lie algebra  $\overline{\mathcal{G}}_{k-1}$  is identified with the additive subgroup of the matrix group  $\mathcal{G}_k$ , whose matrices have the form:

$$\begin{pmatrix} \mathcal{E} & 0\\ \mathcal{A}_k & \mathcal{E} \end{pmatrix},\tag{21}$$

where  $\mathcal{E}$  is the unit block, i.e., the unity of the group  $\mathcal{G}_{k-1}$ . The product and the inversion of such matrices are performed in the following way:

$$\begin{pmatrix} \mathcal{E} & 0\\ \mathcal{A}_k & \mathcal{E} \end{pmatrix} \cdot \begin{pmatrix} \mathcal{E} & 0\\ \mathcal{B}_k & \mathcal{E} \end{pmatrix} = \begin{pmatrix} \mathcal{E} & 0\\ \mathcal{A}_k + \mathcal{B}_k & \mathcal{E} \end{pmatrix},$$
$$\begin{pmatrix} \mathcal{E} & 0\\ \mathcal{A}_k & \mathcal{E} \end{pmatrix}^{-1} = \begin{pmatrix} \mathcal{E} & 0\\ -\mathcal{A}_k & \mathcal{E} \end{pmatrix}.$$

All these matrices generate within the group  $\mathcal{G}_k$  a normal divisor,

$$\begin{pmatrix} \mathcal{A} & 0 \\ \mathcal{A}_k & \mathcal{A} \end{pmatrix} \cdot \begin{pmatrix} \mathcal{E} & 0 \\ \mathcal{B}_k & \mathcal{E} \end{pmatrix} \cdot \begin{pmatrix} \mathcal{A} & 0 \\ \mathcal{A}_k & \mathcal{A} \end{pmatrix}^{-1} = \begin{pmatrix} \mathcal{E} & 0 \\ \mathcal{A}\mathcal{B}_k \mathcal{A}^{-1} & \mathcal{E} \end{pmatrix}.$$

An inner automorphism of the group  $\mathcal{G}_k$  leads to the transformation of the block

$$\mathcal{B}_k \quad \rightsquigarrow \quad \widetilde{\mathcal{B}}_k = -\mathcal{A}\mathcal{B}_k\mathcal{A}^{-1}.$$

Under such a transformation, the spectrum of the matrix  $\mathcal{B}_k$  is preserved. The invariants will be the eigenvalues of this matrix and the corresponding symmetric polynomials, which are coefficients in the Hamilton-Cayley formula.

Exercise 21: Lie algebra of the Lie group. Show that the Lie algebra of an arbitrary Lie group G may be regarded as an additive subgroup and a normal divisor of the tangent group TG. Describe the cosets of this normal divisor and the corresponding quotient group of the group TG.

**Exercise 22: structure constants iterated.** The structure constants of a Lie group G have three indices and can be placed into a spacial matrix a). Prove that the structure constants of the tangent groups TG,  $T^2G$  and  $T^3G$  can be similarly put into a spacial matrices of type b), c) and d), respectively.

## 2. Tangent bundles and osculators

**2.1.** Levels and sector-forms. The tangent functor T iterated k times associates to a smooth manifold M its k-fold tangent bundle  $T^kM$  (the k-th level of M) and associates to a smooth map  $\varphi : M_1 \to M_2$  the graded morphism  $T^k\varphi: T^kM_1 \to T^kM_2$ , the k-th derivative of  $\varphi$ . The level  $T^kM$  has a multiple vector bundle structure with k projections onto  $T^{k-1}M$ :

$$\rho_s \doteq T^{k-s} \pi_s : T^k M \to T^{k-1} M, \quad s = 1, 2, \dots, k,$$

where  $\pi_s$  is the natural projection  $T^s M \to T^{s-1} M$ .

Local coordinates in neighborhoods

$$T^{s}U \subset T^{s}M, \ s = 1, 2, \dots, k, \text{ where } T^{s-1}U = \pi_{s}(T^{s}U),$$

are determined automatically by those in the neighborhood  $U \subset M$ , the quantities  $(u^i)$  being regarded either as coordinate functions on U or as the coordinate components of the point  $u \in U$ :

 $U: \quad (u^{i}), \ i = 1, 2, \dots, n = \dim M, \\ TU: \quad (u^{i}, u^{i}_{1}), \quad \text{with} \quad u^{i} \doteq u^{i} \circ \pi_{1}, \ u^{i}_{1} \doteq du^{i}, \\ T^{2}U: \quad (u^{i}, u^{i}_{1}, u^{i}_{2}, u^{i}_{12}), \\ \vdots \text{ith} \ u^{i} \doteq u^{i} \circ \pi_{2} = u^{i} \circ \pi_{2} \quad u^{i} \circ \pi_{2} \quad u^{i} \doteq d(u^{i} \circ \pi_{2}) \quad u^{i} = u^{i} \circ \pi_{2} \quad u^$ 

with  $u^{i} \doteq u^{i} \circ \pi_{1}\pi_{2}$ ,  $u^{i}_{1} \doteq du^{i} \circ \pi_{2}$ ,  $u^{i}_{2} \doteq d(u^{i} \circ \pi_{1})$ ,  $u^{i}_{12} \doteq d(du^{i})$ , etc.

We set up the following convention: to introduce coordinates on  $T^kU$ , we take the coordinates on  $T^{k-1}U$  and repeat them with an additional index k, so that a tangent vector is preceded by its point of origin. This indexing is convenient since at present the symbols with index s become fiber coordinates for the projection  $\rho_s$ ,  $s = 1, 2, \ldots, k$ .

Thus, for example, under the projections  $\rho_s: T^3U \to T^2U$ , s = 1, 2, 3, the coordinates with indices 1, 2 and 3 are each suppressed in turn:

$$\begin{array}{cccc} (u^{i} \, u_{1}^{i} \, u_{2}^{i} \, u_{12}^{i} \, u_{3}^{i} \, u_{13}^{i} \, u_{23}^{i} \, u_{123}^{i}) \\ \rho_{1} \swarrow & \rho_{2} \downarrow & \searrow \rho_{3} \\ (u^{i} \, u_{2}^{i} \, u_{3}^{i} \, u_{23}^{i}) & (u^{i} \, u_{1}^{i} \, u_{3}^{i} \, u_{13}^{i}) & (u^{i} \, u_{1}^{i} \, u_{2}^{i} \, u_{12}^{i}). \end{array}$$

The level  $T^k M$  is a smooth manifold of dimension  $2^k n$  and admits an important subspace of dimension (k+1)n called the *osculating bundle* of M (briefly – *osculator*) of order k-1 and denoted by  $\operatorname{Osc}^{k-1} M$ . The bundle  $\operatorname{Osc}^{k-1} M$  is determined by the equality of the projections

$$\rho_1=\rho_2=\ldots=\rho_k,$$

meaning that an element of  $T^k M$  belongs to the bundle  $\operatorname{Osc}^{k-1} M$  precisely when all its k projections into  $T^{k-1}M$  coincide. In this case all coordinates with the same number of lower indices coincide. For example, the first bundle  $\operatorname{Osc} M$  is determined in  $T^2 U \subset T^2 M$  by the equation  $u_1^i = u_2^i$ , and the second bundle  $\operatorname{Osc}^2 M$  is determined in  $T^3 U \subset T^3 M$  by  $u_1^i = u_2^i = u_3^i$ ,  $u_{12}^i = u_{13}^i = u_{23}^i$ , etc. The coordinates in  $\operatorname{Osc}^{k-1} M$  will be denoted by the derivatives of the coordinate functions on U, that is  $(u^i, du^i, d^2 u^i, \ldots, d^k u^i)$ .

The immersion  $\zeta$ :  $\operatorname{Osc} M \hookrightarrow T^2 M$  and its derivative  $T\zeta$  are determined in coordinates by matrix formulas:

$$\begin{pmatrix} u^i \\ u^i_1 \\ u^i_2 \\ u^i_{12} \end{pmatrix} \circ \zeta = \begin{pmatrix} u^i \\ du^i \\ du^i \\ d^2u^i \end{pmatrix}, \quad \begin{pmatrix} u^i_3 \\ u^i_{13} \\ u^i_{23} \\ u^i_{123} \end{pmatrix} \circ T\zeta = \begin{pmatrix} du^i \\ d^2u^i \\ d^2u^i \\ d^3u^i \end{pmatrix},$$

$$T\zeta\left(\frac{\partial}{\partial u^{i}}, \ \frac{\partial}{\partial (du^{i})}, \ \frac{\partial}{\partial (d^{2}u^{i})}\right) = \left(\frac{\partial}{\partial u^{i}}, \ \frac{\partial}{\partial u^{i}_{1}} + \frac{\partial}{\partial u^{i}_{2}}, \ \frac{\partial}{\partial u^{i}_{12}}\right)$$

The fibres of the bundle  $\operatorname{Osc} M$  are the integral manifolds of the distribution

$$\langle \partial_i^1 + \partial_i^2, \partial_i^{12} \rangle, \quad \text{with} \quad \partial_i^1 + \partial_i^2 \doteq \frac{\partial}{\partial u_1^i} + \frac{\partial}{\partial u_2^i}, \quad \partial_i^{12} \doteq \frac{\partial}{\partial u_{12}^i}.$$

The functions  $(u_1^i - u_2^i)$  vanish on  $\operatorname{Osc} M$ .

Historically, osculating bundles were introduced under various names long before the bundles  $T^k M$ . The systematic study which was initiated 60 years ago by works of V. Vagner [2] has been culminated in recent times in Miron-Atanasiu theory [3]. Meanwhile, the theme of levels  $T^k M$  remained unjustly neglected for the obvious reason that the multiple fibre bundle structure demands a whole new understanding and new approach (see [1, 4–6]). Attempts such as [7] and the so-called synthetic formulation of  $T^k M$  [8] made progress in that direction.

While an infinitesimal displacement of the point  $u \in M$  is determined by a tangent vector  $u_1$  to M, an infinitesimal displacement of the element  $(u, u_1) \in TM$  is determined by the quantities  $(u_2, u_{12})$ , representing a tangent vector to TM, etc. This interpretation of the elements of  $T^kM$  allows us to develop the theory of higher order motion. Clearly, the future belongs to these bundles.

White considers on the level  $T^k M$  or on a k-multiple vector bundle certain sectorforms which are functions simultaneously linear on the fibres of all k projections (see [7]). In particular, the sector-forms on  $T^2U$  and  $T^3U$  can be written as

$$\begin{split} \Phi &= \varphi_{ij} u_1^i u_2^j + \varphi_i u_{12}^i, \\ \Psi &= \psi_{ijk} u_1^i u_2^j u_3^k + \psi_{ij}^1 u_1^i u_{23}^j + \psi_{ij}^2 u_2^i u_{13}^j + \psi_{ij}^3 u_3^i u_{12}^j + \psi_i u_{123}^i, \end{split}$$

with coefficients in U. For example, in each term of  $\Psi$  the index 1 (or 2 or 3 respectively) appears exactly once. This means that the function  $\Psi$  is linear on the fibres of  $\rho_1$  (and  $\rho_2$  and  $\rho_3$ ).

Any scalar function can be lifted from the level  $T^{k-1}M$  to the level  $T^kM$  by k different projections  $\rho_s : T^kM \to T^{k-1}M$ . For example, for the sector-form  $\Phi$  (see above) there are three possibilities of lifting to  $T^3M$ :

$$\Phi \circ \rho_1 = \varphi_{ij} u_2^i u_3^j + \varphi_i u_{23}^i, \quad \Phi \circ \rho_2 = \varphi_{ij} u_1^i u_3^j + \varphi_i u_{13}^i, \quad \Phi \circ \rho_3 = \varphi_{ij} u_1^i u_2^j + \varphi_i u_{12}^i.$$

**Proposition 1.** Every exterior k-form can be regarded as a sector-form in the sense of White, a scalar function on  $T^kM$  that is constant on the fibres of  $\operatorname{Osc}^{k-1}M$ .

**Proof.** The sector-form  $\Phi$  is constant on  $\operatorname{Osc} M$  if and only if its derivatives vanish on  $\operatorname{Osc} M$ . Thus

$$\begin{split} \Phi &= \varphi_{ij} u_1^i u_2^j + \varphi_i u_{12}^i \quad \Rightarrow \\ & (\partial_i^1 + \partial_i^2) \Phi = \varphi_{ij} u_2^j + \varphi_{ji} u_1^j = (\varphi_{ij} + \varphi_{ji}) u_1^j - \varphi_{ij} (u_1^j - u_2^j), \\ & \partial_i^{12} \Phi = \varphi_i \quad \Rightarrow \quad \varphi_{(ij)} = 0, \quad \varphi_i = 0 \,. \end{split}$$

If  $\Phi$  is an antisymmetric bilinear form then it can be expressed in the coordinates  $(u^i, du^i)$  as a 2-form  $\Phi = \varphi_{[ij]} du^i \wedge du^j$ . Thus the sector-form  $\Phi$  is constant on  $\operatorname{Osc} M$  if and only if it is a Cartan 2-form.

If k = 3 the fibres  $\operatorname{Osc}^2 M$  of dimension 3n are the integral manifolds of the distribution

$$\langle \partial_i^1 + \partial_i^2 + \partial_i^3, \ \partial_i^{23} + \partial_i^{13} + \partial_i^{12}, \ \partial_i^{123} \rangle.$$

For the sector-form  $\Psi$  (see above) we have

$$\begin{split} \Psi &= \psi_{ijk} u_1^i u_2^j u_3^k + \psi_{ij}^1 u_1^i u_{23}^j + \psi_{ij}^2 u_2^i u_{13}^j + \psi_{ij}^3 u_3^i u_{12}^j + \psi_i u_{123}^i \Rightarrow \\ (\partial_i^1 + \partial_i^2 + \partial_i^3) \Psi &= \psi_{ijk} u_2^j u_3^k + \psi_{jik} u_1^j u_3^k + \psi_{jki} u_1^j u_2^k + \psi_{ij}^1 u_{23}^j + \psi_{ij}^2 u_{13}^j + \psi_{ij}^3 u_{12}^j , \\ (\partial_i^{23} + \partial_i^{13} + \partial_i^{12}) \Psi &= \psi_{ji} u_1^j + \psi_{ji}^2 u_2^j + \psi_{ji}^3 u_3^j , \\ \partial_i^{123} \Psi &= \psi_i . \end{split}$$

The derivatives vanish on the fibres  $Osc^2M$  when the following conditions hold:

$$\varphi_{(ijk)} = 0, \quad \psi_{ij}^1 + \psi_{ij}^2 + \psi_{ij}^3 = 0, \quad \psi_i = 0.$$

These conditions are necessary and sufficient for the sector-form  $\Psi$  to be constant on  $\operatorname{Osc}^2 M$ , but not for  $\Psi$  to be a Cartan 3-form. However, every 3-form  $\widetilde{\Psi} = \varphi_{ijk} du^i \wedge du^j \wedge du^k$  can be regarded as a homogeneous sector-form that is constant on  $\operatorname{Osc}^2 M$ .

The argument extends likewise to the cases when k > 3.

White's theory of sector-forms is much more extensive than that of Cartan exterior forms. In particular, exterior differentiation is an operation on the set of sector-forms that are constant on the osculating bundles.

**2.2.** Gauge groups on osculating spaces. The action of the gauge group  $\mathcal{G}_k$  on the k-th level  $T^k M$  extends in a natural way to the osculating bundle  $\operatorname{Osc}^{k-1} M$ . The diagram from below shows how the block-matrix  $4 \times 4$  reduces, for  $u_1 = u_2$ , to a  $3 \times 3$  block-matrix:

$$u_1 = u_2 \quad \Rightarrow \quad \begin{pmatrix} a & 0 & 0 & 0 \\ a_1 & a & 0 & 0 \\ a_2 & 0 & a & 0 \\ a_{12} & a_2 & a_1 & a \end{pmatrix} \quad \rightsquigarrow \quad \begin{pmatrix} a & 0 & 0 \\ da & a & 0 \\ d^2a & da & a \end{pmatrix}.$$

The blocks of the matrix from the right side are generated in the following way:

The action of the gauge group  $\mathcal{G}_2$  on the level  $T^2M$  is obviously transported to the subbundle  $\operatorname{Osc} M \subset T^2M$ . While one passes from  $T^2M$  to  $\operatorname{Osc} M$  by considering

$$(a_1 = a_2, a_{12}) \rightsquigarrow (da, d^2a), \quad (\partial^1 + \partial^2, \partial^{12}) \rightsquigarrow \left(\frac{\partial}{\partial(du)}, \frac{\partial}{\partial(d^2u)}\right),$$

the transformation of the natural basis on  $T^2M$  is transported to the transformation of the natural basis on  $\operatorname{Osc} M$ :

$$(\partial \ \partial^1 \ \partial^2 \ \partial^{12}) \cdot \begin{pmatrix} a & 0 & 0 & 0 \\ a_1 & a & 0 & 0 \\ a_2 & 0 & a & 0 \\ a_{12} & a_2 & a_1 & a \end{pmatrix} \rightsquigarrow \begin{pmatrix} \frac{\partial}{\partial u} \ \frac{\partial}{\partial (du)} \ \frac{\partial}{\partial (d^2u)} \end{pmatrix} \cdot \begin{pmatrix} a & 0 & 0 \\ da & a & 0 \\ d^2a & da & a \end{pmatrix}.$$

In the general case, the action of the group  $\mathcal{G}_k$  on the level  $T^k M$  extends in a similar way to the subbundle  $\operatorname{Osc}^{k-1} M$ .

#### Резюме

М. Рахула, В. Балан. Касательные расслоения и калибровочные группы.

Дифференциалы  $T^k a$   $(k \ge 1)$  диффеоморфизма a гладкого многообразия M индуцируют в слоях расслоений  $T^k M$ , то есть в соответствующих касательных пространствах, линейные преобразования, заключающие в себе действие калибровочной группы  $\mathcal{G}_k$ . Это действие естественным образом распространяется на соприкасающиеся подрасслоения  $\operatorname{Osc}^{k-1} M \subset T^k M$ .

Ключевые слова: диффеоморфизм гладкого многообразия, пространство расслоения, действие калибровочной группы.

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