# TANGENT BUNDLES AND GAUGE GROUPS 

M. Rahula, V. Balan


#### Abstract

The differentials $T^{k} a(k \geq 1)$ of a diffeomorphism $a$ of a smooth manifold $M$ induce in the fibers of the fiber bundles $T^{k} M$, i.e., in the corresponding tangent spaces, linear transformations, which embody the action of the gauge group $\mathcal{G}_{k}$. This action extends in a natural way to the osculating subbundles $\mathrm{Osc}^{k-1} M \subset T^{k} M$.


Key words: diffeomorphism of a smooth manifold, fiber bundles, action of the gauge group.

## Introduction

The differential group $\mathcal{G}$ of a smooth manifold $M$ induces in the tangent bundle $T^{k} M$ an action of the group of $k$-jets of transformations. More specifically, if $a$ is a diffeomorphism of the manifold $M$, then its $k$-th differential $T^{k} a$ is a transformation of the level $T^{k} M$. Then the level $T^{k} M$ may be regarded as a homogeneous space $J^{k} / H_{k}$, where $J^{k}$ is the group of $k$-jets of transformations and $H_{k}$ is the stabilizer of an element $u_{(u)} \in T^{k} M$. The gauge group $\mathcal{G}_{k}$ is defined as a certain subgroup of the linear group $G L\left(2^{k} n, \mathbb{R}\right)$, where $n=\operatorname{dim} M$ which is isomorphic to the stabilizer $H_{k}$. The action of the group $\mathcal{G}$ extends to the osculating subbundle $\mathrm{Osc}^{k-1} M \subset T^{k} M$.

The paper contains all the necessary definitions and founds all the previous considerations. Commented examples and groups of derived formulas are presented as exercises.

## 1. Tangent groups

1.1. Leibniz rule. We apply the tangent functor $T$ to the Cartesian product of smooth manifolds $M_{1}$ and $M_{2}$ :

$$
T\left(M_{1} \times M_{2}\right)=\left(T M_{1} \times M_{2}\right) \oplus\left(M_{1} \times T M_{2}\right)
$$

and for the smooth mapping from $M_{1} \times M_{2}$ to some smooth manifold $M$

$$
\lambda: M_{1} \times M_{2} \longrightarrow M:(u, v) \mapsto w=u \cdot v
$$

we define the tangent mapping $T \lambda$. First, by fixing the points $u \in M_{1}$ and $v \in M_{2}$ we define two mappings $\lambda_{u}$ and $\lambda_{v}$ :

$$
\lambda_{u}: M_{2} \rightarrow M: v \mapsto u \cdot v, \quad \lambda_{v}: M_{1} \rightarrow M: u \mapsto u \cdot v
$$

Theorem 1. To the pair of vectors $u_{1} \in T_{u} M_{1}$ and $v_{1} \in T_{v} M_{2}$ the mapping $T \lambda$ associates the vector $w_{1} \in T_{w} M$, and we have

$$
\begin{equation*}
w=u \cdot v \quad \Rightarrow \quad w_{1}=u_{1} \cdot v+u \cdot v_{1} \tag{1}
\end{equation*}
$$

where $u_{1} \cdot v=T \lambda_{v}\left(u_{1}\right)$ and $u \cdot v_{1}=T \lambda_{u}\left(v_{1}\right)$. In short, one can apply to the "product" $w=u \cdot v$ the Leibniz rule.

Proof. We specify that, by means of the tangent maps $T \lambda_{v}$ and $T \lambda_{u}$, two vectors $u_{1} \in T_{u} M_{1}$ and $v_{1} \in T_{v} M_{2}$ are transported from the points $u \in M_{1}$ and $v \in M_{2}$ to the point $w \in M$, where their sum defines the vector $w_{1} \in T_{w} M$. Locally, this is confirmed by the formula:

$$
w^{\rho}=\lambda^{\rho}\left(u^{i}, v^{\alpha}\right) \quad \Rightarrow \quad w_{1}^{\rho}=\frac{\partial \lambda^{\rho}}{\partial u^{i}} u_{1}^{i}+\frac{\partial \lambda^{\rho}}{\partial v^{\alpha}} v_{1}^{\alpha}
$$

where $u^{i}, v^{\alpha}, w^{\rho}$ are the coordinates of the points $u, v, w$ and $u_{1}^{i}, v_{1}^{\alpha}, w_{1}^{\rho}$ are the components of the vectors $u_{1}, v_{1}, w_{1}$ on the neighborhoods $U_{1} \subset M_{1}, U_{2} \subset M_{2}, U \subset M$, $i=1, \ldots, \operatorname{dim} M_{1}, \alpha=1, \ldots, \operatorname{dim} M_{2}, \rho=1, \ldots, \operatorname{dim} M$.

Using the Leibniz rule we derive a set of important formulas in coordinate free form.
Exercise 1: action of Leibniz rule. Show that the Leibniz rule can be applied to the "product" of several factors, e.g.,

$$
(u \cdot v \cdot w)_{1}=u_{1} \cdot v \cdot w+u \cdot v_{1} \cdot w+u \cdot v \cdot w_{1}
$$

Exercise 2: prolongation of Leibniz rule. Prove that for the second tangent mapping $T^{2} \lambda$, the following formulas hold true:

$$
\begin{align*}
w=u \cdot v, \quad w_{1} & =u_{1} \cdot v+u \cdot v_{1} \\
w_{2} & =u_{2} \cdot v+u \cdot v_{2}  \tag{2}\\
w_{12} & =u_{12} \cdot v+u_{2} \cdot v_{1}+u_{1} \cdot v_{2}+u \cdot v_{12}
\end{align*}
$$

## Exercise 3: functional equation.

Question: how can one solve the equation $(u \cdot v)_{1}=u_{1} \cdot v+u \cdot v_{1}$ with respect to $u_{1}$ for given $v_{1}$ and $(u \cdot v)_{1}$, or relative to $v_{1}$ for given $u_{1}$ and $(u \cdot v)_{1}$ ? This reminds the method of integration by parts:

$$
\begin{aligned}
& d(u v)=u d v+v d u \quad u v=\int u d v+\int v d u, \quad \text { whence } \\
& \text { either } \int u d v=u v-\int v d u, \quad \text { or } \quad \int v d u=u v-\int u d v .
\end{aligned}
$$

1.2. Coordinate-free story. The rule (1) is easy to use while building tangent groups and further, while studying representations of groups. If we have previously denoted the "product" of elements by a dot, as in (1) and (2), then while denoting the product of group elements, the dot will be omitted.

To a Lie group $G$ with composition rule $\gamma:(a, b) \mapsto c=a b$, we associate the tangent group $T G$, having the composition law $T \gamma$ :

$$
\begin{equation*}
c=a b \quad \Rightarrow \quad c_{1}=a_{1} b+a b_{1} \tag{3}
\end{equation*}
$$

The vectors $a_{1} \in T_{a} G$ and $b_{1} \in T_{b} G$ are transported by means of the right shift $r_{b} \doteq \gamma_{b}$ and of the left shift $l_{a} \doteq \gamma_{a}$, more exactly, by means of the tangent mappings $\operatorname{Tr}_{b}$ and $T l_{a}$, from the points $a$ and $b$ to the point $c$, where the sum $a_{1} b+a b_{1}=T r_{b}\left(a_{1}\right)+T l_{a}\left(b_{1}\right)$ determines the vector $c_{1} \in T_{c} G$. This is the composition law on the tangent group $T G$. The unity of the group $T G$ is the null vector from $T_{e} G$. The inversion for the elements of $T G$ is defined by the rule:

$$
\begin{equation*}
a_{1} \in T_{a} G \quad \rightsquigarrow \quad a_{1}^{-1}=-a^{-1} a_{1} a^{-1} \in T_{a-1} G . \tag{4}
\end{equation*}
$$

Exercise 4: unity and inverse elements. Using (3), confirm the assertion regarding the unity of the group $T G$ and the inversion of the elements (4). The formula (4) is obtained by solving the equation $a_{1} b+a b_{1}=0$ relative to $b_{1}$ for $b=a^{-1}$.

Exercise 5: matrix representation. Prove that the formulas (3) can be represented in matrix form as

$$
\left(\begin{array}{cc}
a & 0  \tag{5}\\
a_{1} & a
\end{array}\right) \cdot\left(\begin{array}{cc}
b & 0 \\
b_{1} & b
\end{array}\right)=\left(\begin{array}{cc}
a & 0 \\
0 & a
\end{array}\right) \cdot\left(\begin{array}{cc}
e & 0 \\
a^{-1} a_{1}+b_{1} b^{-1} & e
\end{array}\right) \cdot\left(\begin{array}{ll}
b & 0 \\
0 & b
\end{array}\right) \cdot
$$

The sum of vectors in $T_{a b} G$ reduces to the sum in $T G_{e}$ :

$$
a_{1} b+a b_{1}=T\left(l_{a} \circ r_{b}\right)\left(a^{-1} a_{1}+b_{1} b^{-1}\right)
$$

Explain the meaning of the equality $a_{1} b+a b_{1}=\left(a_{1} a^{-1}\right) c+c\left(b^{-1} b_{1}\right)$.
Exercise 6: second tangent group. Prove that in the second tangent group $T^{2} G$ the product of elements is defined by the formulas

$$
\begin{align*}
c=a b, \quad c_{1} & =a_{1} b+a b_{1}, \\
c_{2} & =a_{2} b+a b_{2},  \tag{6}\\
c_{12} & =a_{12} b+a_{2} b_{1}+a_{1} b_{2}+a b_{12},
\end{align*}
$$

and the inversion is performed by the rule

$$
\begin{align*}
& \left(a, a_{1}, a_{2}, a_{12}\right)^{-1} \doteq\left(a^{-1}, a_{1}^{-1}, a_{2}^{-1}, a_{12}^{-1}\right), \quad \text { where } \\
& a_{1}^{-1}=-a^{-1} a_{1} a^{-1} \\
& a_{2}^{-1}=-a^{-1} a_{2} a^{-1}  \tag{7}\\
& a_{12}^{-1}=-a^{-1} a_{12} a^{-1}+a^{-1} a_{2} a^{-1} a_{1} a^{-1}+a^{-1} a_{1} a^{-1} a_{2} a^{-1} .
\end{align*}
$$

Exercise 7: classical formulas. Reduce the formulas (6) and (7) to the well known formulas from Analysis:

$$
\begin{aligned}
& (u v)^{\prime}=u^{\prime} v+u v^{\prime}, \quad(u v)^{\prime \prime}=u^{\prime \prime} v+2 u^{\prime} v^{\prime}+u v^{\prime \prime} \\
& \left(\frac{1}{u}\right)^{\prime}=-\frac{u^{\prime}}{u^{2}}, \quad\left(\frac{1}{u}\right)^{\prime \prime}=\frac{-u u^{\prime \prime}+2\left(u^{\prime}\right)^{2}}{u^{3}}
\end{aligned}
$$

Exercise 8: matrix relations. Using (6) and (7) prove the matrix relations:

$$
\begin{gathered}
\left(\begin{array}{cccc}
c & 0 & 0 & 0 \\
c_{1} & c & 0 & 0 \\
c_{2} & 0 & c & 0 \\
c_{12} & c_{2} & c_{1} & c
\end{array}\right)=\left(\begin{array}{cccc}
a & 0 & 0 & 0 \\
a_{1} & a & 0 & 0 \\
a_{2} & 0 & a & 0 \\
a_{12} & a_{2} & a_{1} & a
\end{array}\right) \cdot\left(\begin{array}{ccc}
b & 0 & 0 \\
b_{1} & b & 0 \\
0 \\
b_{2} & 0 & b \\
b_{12} & b_{2} & b_{1} \\
b
\end{array}\right) \\
\left(\begin{array}{cccc}
a & 0 & 0 & 0 \\
a_{1} & a & 0 & 0 \\
a_{2} & 0 & a & 0 \\
a_{12} & a_{2} & a_{1} & a
\end{array}\right)^{-1}=\left(\begin{array}{cccc}
a^{-1} & 0 & 0 & 0 \\
a_{1}^{-1} & a^{-1} & 0 & 0 \\
a_{2}^{-1} & 0 & a^{-1} & 0 \\
a_{12}^{-1} & a_{2}^{-1} & a_{1}^{-1} & a^{-1}
\end{array}\right)
\end{gathered}
$$

Which endomorphism is involved here?
Exercise 9: logarithmic derivatives. Following the example of (5), represent the product of elements of the group $T^{2} G$ in the form:

$$
\left(\begin{array}{cccc}
a & 0 & 0 & 0 \\
a_{1} & a & 0 & 0 \\
0 & 0 & a & 0 \\
0 & 0 & a_{1} & a
\end{array}\right) \cdot\left(\begin{array}{cccc}
e & 0 & 0 & 0 \\
0 & e & 0 & 0 \\
a^{-1} a_{2}+b_{2} b^{-1} & 0 & e & 0 \\
\left(a^{-1} a_{2}+b_{2} b^{-1}\right)_{1} & a^{-1} a_{2}+b_{2} b^{-1} & 0 & e
\end{array}\right) \cdot\left(\begin{array}{cccc}
b & 0 & 0 & 0 \\
b_{1} & b & 0 & 0 \\
0 & 0 & b & 0 \\
0 & 0 & b_{1} & b
\end{array}\right),
$$

where $\left(a^{-1} a_{2}+b_{2} b^{-1}\right)_{1}=a^{-1} a_{12}-a^{-1} a_{1} a^{-1} a_{2}+b_{12} b^{-1}-b_{2} b^{-1} b_{1} b^{-1}$. In this generalization, both logarithmic derivatives $(\ln (u v))^{\prime}$ and $(\ln (u v))^{\prime \prime}$, are present.

At the unity $e \in G$ we fix the tangent vector $e_{1} \in T_{e} G$. This vector is displaced by left shifts $l_{a}$ over the group $G$ to produce the left-invariant vector field $a e_{1}$ and by right shifts $r_{a}$, to produce the right-invariant vector field $e_{1} a$. If at the unit $e \in G$ we provide a frame, i.e., a basis of the space $T_{e} G$, then, in this way, two frame fields are defined on $G$ (one right-invariant and left-invariant). The transition from one frame to another, at the point $a \in G$, is defined by some matrix $A(a)$, which is an element of the linear group $G L=G L(\operatorname{dim} G, \mathbb{R})$. By this way, we define a homomorphism of the group $G$ into the linear group $G L$ :

$$
\begin{equation*}
G \rightarrow G L: a \mapsto A(a) . \tag{8}
\end{equation*}
$$

Exercise 10: right/left shifts and inner automorphisms. Show that an 1-parametric subgroup $a_{t}$ of the group $G$ defines in the group $G$ three flows corresponding to right shifts, left shifts and inner automorphisms

$$
r_{a_{t}}=\exp t X, \quad l_{a_{t}}=\exp t \widetilde{X}, \quad A_{a_{t}}=l_{a_{t}} \circ r_{a_{t}}^{-1}=\exp t(\tilde{X}-X)
$$

and, accordingly, the left-invariant operator $X$, the right-invariant operator $\widetilde{X}$ and the adjoint representation operator $Y=\widetilde{X}-X$. Prove this, using the formulas

$$
X f=\left(f \circ r_{a_{t}}\right)_{t=0}^{\prime}, \quad \widetilde{X} f=\left(f \circ l_{a_{t}}\right)_{t=0}^{\prime}, \quad Y f=\left(f \circ A_{a_{t}}\right)_{t=0}^{\prime}
$$

where $f$ is an arbitrary smooth function on $G$, taking into consideration that left shifts commute with right shifts.
1.3. Elements of representation theory. We consider a differentiable manifold, which we shall call representation space for the group $G$, or, in the following, simply space. A smooth mapping

$$
\lambda: M \times G \longrightarrow M:(u, a) \mapsto v=u \cdot a
$$

defines an action of the group $G$ on the space $M$, if all the mappings

$$
\lambda_{a}: M \rightarrow M u \mapsto u \cdot a, \quad \forall a \in G
$$

are transformations (diffeomorphisms) of the space $M$, and the mapping $a \mapsto \lambda_{a}$ is a homomorphism of the group $G$ into the group of transformations of the space $M$. The homomorphism $a \mapsto \lambda_{a}$ is understood either in the sense of the equality $\lambda_{a b}=\lambda_{a} \circ \lambda_{b}$ or in the sense of the equality $\lambda_{a b}=\lambda_{b} \circ \lambda_{a}$. In the first case we say that the action of the group $G$ on the space $M$ is left-sided, and in the second case it is right-sided. By writing $v=a \cdot u$ we have in view a left action, while by $v=u \cdot a$ - a right action:

$$
\begin{aligned}
& v=a \cdot u \quad \rightsquigarrow \quad(a b) \cdot u=a \cdot(b \cdot u), \\
& v=u \cdot a \quad \rightsquigarrow \quad u \cdot(a b)=(u \cdot a) \cdot b .
\end{aligned}
$$

The next formulas correspond to the right-sided action of the group $G$ on the space $M$.
The kernel of the homomorphism $a \mapsto \lambda_{a}$ is called the stabilizer subgroup of the group $G$. In the case of an effective action, the non-effectiveness kernel is trivial; it consists of the unity $e \in G$, and the mapping $a \mapsto \lambda_{a}$ is injective.

For a fixed point $u \in M$, the mapping

$$
\lambda_{u}: G \rightarrow M: a \mapsto u \cdot a
$$

defines in the space $M$ the orbit of this point. The whole space $M$ is fibered into orbits $\lambda_{u}(G)$. When $\lambda_{u}(G)=M$, i.e., when the space $M$ is the only orbit of the group $G$, we say that the action of the Lie group $G$ on the space $M$ is transitive. In such a case, $M$ is called homogeneous group space. If moreover, the dimensions of $G$ and $M$ are equal, then the action of $G$ on $M$ is simply transitive and such an action defines an exact representation of the group $G$.

Equation 11: action by right/left shifts and inner automorphisms. Show that the actions of the group $G$ on itself, provided by left and right shifts, are simply transitive actions. Prove that the actions provided by inner automorphisms are non-transitive.

The tangent map of the mapping $\lambda$, i.e., $T \lambda$, defines a representation of the tangent group $T G$ on the first level $T M$,

$$
T \lambda: T M \times T G \rightarrow T M:\left(u_{1}, a_{1}\right) \mapsto v_{1}=u_{1} \cdot a+u \cdot a_{1}
$$

Formula (1) looks similarly:

$$
\begin{equation*}
v=u \cdot a \quad \rightsquigarrow \quad v_{1}=u_{1} \cdot a+u \cdot a_{1} . \tag{9}
\end{equation*}
$$

We remark two particular cases:
for $a_{1}=0$ we define the action of the group $G$ on the level $T M$,

$$
a_{1}=0 \quad \Rightarrow \quad u_{1} \mapsto v_{1}=u_{1} \cdot a ;
$$

for $u_{1}=0$ we define the action of the tangent group $T G$ on the space $M$,

$$
u_{1}=0 \quad \Rightarrow \quad e_{1}=a^{-1} a_{1} \quad \mapsto \quad v_{1}=u \cdot a_{1}=v \cdot a^{-1} a_{1}=v \cdot e_{1}
$$

The formula

$$
\begin{equation*}
v_{1}=v \cdot a^{-1} a_{1} \tag{10}
\end{equation*}
$$

is the fundamental formula of the theory of Lie group representations.
In fact, to the vector $e_{1}=a^{-1} a_{1} \in T_{e} G$ (which is an element of the Lie algebra $g$ ) at the point $v \in M$ we associate some vector $v_{1} \in T_{v} M$, and since $v$ is an arbitrary point, it defines a vector field on the space $M$. This vector field is called the operator of the group $G$, or simply group operator ${ }^{1}$. Depending on the choice of the vector $e_{1} \in T_{e} G$, in the space $M$ we have an infinite set of group operators, and all of them, as vector fields, are tangent to the corresponding orbits.

For $v_{1}=0$ the equality (10) provides the equation $v \cdot a^{-1} a_{1}=0$, or $v \cdot e_{1}=0$, which determines in the space $T_{e} G$ those directions $e_{1}$, along which the point $v \in M$ remains fixed. We define on the group $G$ a Pfaff system, and its integral surface (solution), which contains the point $e \in G$, is a subgroup $H_{v} \subset G$ called the stationary subgroup or the stabilizer of the point $v$.

In coordinates $\left(v^{\alpha}\right)$ on the neighborhood $U \subset M$ of the point $v \in M$, Eq. (10) is written as a system $d v^{\alpha}=\xi_{i}^{\alpha} \omega^{i}$, where $\omega^{i}$ are the forms of the left-invariant coframe on the group $G$. There appears a matrix ${ }^{2} \xi=\left(\xi_{i}^{\alpha}\right)$, which determines a system of forms $\vartheta^{\alpha}$ on the group $G$, and in the space $M$, a system of basic operators $X_{i}$ :

$$
\vartheta^{\alpha}=\xi_{i}^{\alpha} \omega^{i}, \quad X_{i}=\xi_{i}^{\alpha} \frac{\partial}{\partial v^{\alpha}} .
$$

The number of operators $X_{i}$ is equal to the dimension of $G$, and the number of forms $\vartheta^{\alpha}$ is the dimension of $M$. Operators $X_{i}$ and forms $\vartheta^{\alpha}$ are not necessarily linearly independent. The Pfaff system $\xi_{i}^{\alpha} \omega^{i}=0$ for a fixed point $v \in M$ is completely integrable and defines the stabilizer $H_{v} \subset G$.

Exercise 12: vision from the classical theory. Show that the system $\xi_{i}^{\alpha} \omega^{i}=0$ is the coordinate form of the equation $v \cdot a^{-1} a_{1}=0$.
1.4. Adjoint representation. In the groups $G$ and $T G$ we define the action by left shifts:

$$
\begin{align*}
l_{a} & : b \mapsto c=a b, \\
T l_{a_{1}} & : b_{1} \mapsto c_{1}=\left(a_{1} a^{-1}\right) c+a b_{1}, \\
& c_{1}=\left(a_{1} a^{-1}\right) c, \tag{11}
\end{align*}
$$

[^0]right shifts:
\[

$$
\begin{align*}
r_{a} & : b \mapsto c=b a, \\
\operatorname{Tr}_{a_{1}} & : b_{1} \mapsto c_{1}=b_{1} a+c\left(a^{-1} a_{1}\right), \\
& c_{1}=c\left(a^{-1} a_{1}\right), \tag{12}
\end{align*}
$$
\]

and inner automorphisms:

$$
\begin{align*}
& A_{a}: b \mapsto c=a b a^{-1}, \\
& T A_{a_{1}}: b_{1} \mapsto c_{1}=\left(a_{1} a^{-1}\right) c-c\left(a_{1} a^{-1}\right)+a b_{1} a^{-1}, \\
& c_{1}=\left(a_{1} a^{-1}\right) c-c\left(a_{1} a^{-1}\right) . \tag{13}
\end{align*}
$$

The basic formula (10) is rewritten, for $b_{1}=0$, in the forms (11), (12) and (13), respectively.

Inner automorphisms are directly related to higher order movements.
Hence, when in the spaces $A$ and $B$ there take place the transformations $a$ and $b$, the mapping $\varphi: A \rightarrow B$ is brought into the mapping $\widetilde{\varphi}: A \rightarrow B$. This is shown by the diagram:


If we set here $A=B, a=b$ and if $\varphi$ is a diffeomorphism, i.e., $\varphi$ is a transformation of the space $A$, then this diagram describes the transformation of the mapping $\varphi$, subject to the influence of the transformation $a$ :


The transformation $\varphi$ is subject to the inner automorphism.
Exercise 13: higher order transformations. The transformation of order $2 \varphi \rightsquigarrow \widetilde{\varphi}$ is described by the 2 -dimensional diagram (1.14). Show that the transformation of order 3, i.e., a transformation of transformation $\varphi \rightsquigarrow \widetilde{\varphi}$, is described by a 3-dimensional diagram and the transformation of order $k$ is described by a corresponding $k$-dimensional diagram.

If the arrow $a$ in diagram (14) is assumed to represent the 1-parametric group $a_{t}$ of transformations of the space $A$, or in brief, the flow $a_{t}$, then we see how, to a change of the parameter $t$ (of time), it corresponds to a change of the mapping $\varphi_{t}=a_{t} \varphi a_{t}^{-1}$. We can talk then about a 1-parametric family of mappings $\varphi_{t}$ in the field $a_{t}$.

If the arrow $\varphi$ in diagram (14) is regarded as a 1-parametric group of transformations $b_{\tau}$ of the space $A$ (the flow $b_{\tau}$ ), then we can see how this flow changes under the transformation $a$, i.e., $b_{\tau} \rightsquigarrow \widetilde{b}_{\tau}=a b_{\tau} a^{-1}$.

Exercise 14: transformation of the flow. Show that if $b_{\tau}$ is the flow of the vector field $Y$ and $\widetilde{b}_{\tau}$ is the flow of the field $\widetilde{Y}$, then $\widetilde{Y}=T a Y$, and the tangent mapping $T a$ acts on the field $Y$ :

$$
b_{\tau}=\exp \tau Y \quad \rightsquigarrow \quad \widetilde{b}_{\tau}=a b_{\tau} a^{-1}=\exp \tau \tilde{Y}, \quad Y \rightsquigarrow \tilde{Y}=T a Y
$$

Exercise 15: interaction of vector fields. Let $X$ and $Y$ be two vector fields. The flows of these fields $a_{t}=\exp t X$ and $b_{\tau}=\exp \tau Y$ interact according to the scheme:

$$
\begin{array}{lll}
b_{\tau} \rightsquigarrow a_{t} b_{\tau} a_{-t}, & Y \rightsquigarrow T a_{t} Y, \\
a_{t} \rightsquigarrow b_{\tau} a_{t} b_{-\tau}, & X \rightsquigarrow T b_{\tau} X .
\end{array}
$$

Using derivatives of the function $f$

$$
X f=\left(f \circ a_{t}\right)_{t=0}^{\prime} \quad \text { and } \quad Y f=\left(f \circ b_{\tau}\right)_{\tau=0}^{\prime}
$$

perform the differentiation (the parameters $t$ or $\tau$ from above the arrow mean differentiation relative to $t$ for $t \rightarrow 0$ or to $\tau$, for $\tau \rightarrow 0$ ):

$$
\begin{array}{lll}
f \circ\left(a_{t} b_{\tau} a_{t}^{-1}\right) & \xrightarrow{t}(X f) \circ b_{\tau}-X\left(f \circ b_{\tau}\right) \quad \xrightarrow{\tau}(Y X-X Y) f, \\
f \circ\left(b_{\tau} a_{t} b_{\tau}^{-1}\right) & \xrightarrow{\tau}(Y f) \circ a_{t}-Y\left(f \circ a_{t}\right) \quad \xrightarrow{t}(X Y-Y X) f .
\end{array}
$$

Check the validity of the relation $\left(T a_{t} Y\right)_{t=0}^{\prime}=-\left(T b_{\tau} X\right)_{\tau=0}^{\prime}$ and establish a connection with the brackets $[X, Y]=X Y-Y X$.

If in one flow the points move along trajectories and under the influence of the other flow, this movement is transformed, and then the movement of the movement takes place, or a second-order movement. Under the influence of a third flow, the movement of second order changes its shape, then the movement of third order occurs, etc. In the infinitesimal approach this reduces to the iterations

$$
a_{t} \rightsquigarrow T a_{t} \rightsquigarrow T^{2} a_{t} \rightsquigarrow \ldots
$$

and to the corresponding vector fields on the levels

$$
\begin{equation*}
X \rightsquigarrow \stackrel{(1)}{X} \rightsquigarrow \stackrel{(2)}{X} \rightsquigarrow \ldots, \quad T^{k} a_{t}=\exp t \stackrel{(k)}{X}, \quad k=0,1,2, \ldots \tag{15}
\end{equation*}
$$

In this way, the flow $T^{k} a_{t}$ induces a movement of order $k$.
1.5. Gauge groups. Let us fix on each level a point

$$
u_{(k)} \in T^{k} M, \quad \text { such that } \quad \pi_{k}\left(u_{(k)}\right)=u_{(k-1)}, \quad k=0,1,2, \ldots
$$

In the neighborhood $T^{k} U \subset T^{k} M$ these points are defined by their coordinates:

$$
\begin{aligned}
& U: u_{(0)} \rightsquigarrow\left(u^{i}\right), \\
& T U: u_{(1)} \\
& \rightsquigarrow\left(u^{i}, u_{1}^{i}\right), \\
& T^{2} U: u_{(2)} \rightsquigarrow\left(u^{i}, u_{1}^{i}, u_{2}^{i}, u_{12}^{i}\right), \\
& T^{3} U: u_{(3)} \rightsquigarrow\left(u^{i}, u_{1}^{i}, u_{2}^{i}, u_{12}^{i}, u_{3}^{i}, u_{13}^{i}, u_{23}^{i}, u_{123}^{i}\right), \\
& \ldots
\end{aligned}
$$

Transformation of coordinates in the neighborhood $U \subset M$

$$
u^{i} \rightsquigarrow \widetilde{u}^{i} \circ a=a^{i}
$$

induces a change of coordinates in each neighborhood $T^{k} U$ :

$$
\left(u^{i}, u_{1}^{i}, u_{2}^{i}, u_{12}^{i}, \ldots\right) \rightsquigarrow\left(\widetilde{u}^{i}, \widetilde{u}_{1}^{i}, \widetilde{u}_{2}^{i}, \widetilde{u}_{12}^{i}, \ldots\right)=\left(a^{i}, a_{1}^{i}, a_{2}^{i}, a_{12}^{i}, \ldots\right) .
$$

Namely, if these transformations of coordinates in the neighborhood $T U$ are defined by the system

$$
\left\{\begin{array}{l}
\widetilde{u}^{i}=a^{i}, \\
\widetilde{u}_{1}^{i}=a_{1}^{i} \doteq a_{j}^{i} u_{1}^{j},
\end{array}\right.
$$

with the Jacobian block-matrix

$$
\left(\begin{array}{cc}
a_{j}^{i} & 0 \\
\left(a_{j}^{i}\right)_{1} & a_{j}^{i}
\end{array}\right), \quad \text { where } \quad a_{j}^{i}=\frac{\partial a^{i}}{\partial u^{j}}, \quad a_{j k}^{i}=\frac{\partial^{2} a^{i}}{\partial u^{j} \partial u^{k}}, \quad\left(a_{j}^{i}\right)_{1} \doteq a_{j k}^{i} u_{1}^{k},
$$

then the transformation of coordinates in the neighborhood $T^{2} U$ are defined by the system

$$
\left\{\begin{array}{l}
\widetilde{u}^{i}=a^{i} \\
\widetilde{u}_{1}^{i}=a_{1}^{i} \doteq a_{j}^{i} u_{1}^{j} \\
\widetilde{u}_{2}^{i}=a_{2}^{i} \doteq a_{j}^{i} u_{2}^{j} \\
\widetilde{u}_{12}^{i}=a_{12}^{i} \doteq a_{j k}^{i} u_{1}^{j} u_{2}^{k}+a_{j}^{i} u_{12}^{j}
\end{array}\right.
$$

with the Jacobian block-matrix

$$
\left(\begin{array}{cccc}
a_{j}^{i} & 0 & 0 & 0 \\
\left(a_{j}^{i}\right)_{1} & a_{j}^{i} & 0 & 0 \\
\left(a_{j}^{i}\right)_{2} & 0 & a_{j}^{i} & 0 \\
\left(a_{j}^{i}\right)_{12} & \left(a_{j}^{i}\right)_{2} & \left(a_{j}^{i}\right)_{1} & a_{j}^{i}
\end{array}\right), \quad \text { where } \quad\left\{\begin{array}{l}
\left(a_{j}^{i}\right)_{1} \doteq a_{j k}^{i} u_{1}^{k}, \\
\left(a_{j}^{i}\right)_{2} \doteq a_{j k}^{i} u_{2}^{k}, \\
\left(a_{j}^{i}\right)_{12} \doteq a_{j k l}^{i} u_{1}^{k} u_{2}^{l}+a_{j k}^{i} u_{12}^{k},
\end{array}\right.
$$

etc.
When performing a lift from one level to another, $U \rightsquigarrow T U \rightsquigarrow T^{2} U \rightsquigarrow \ldots$, the Jacobian matrix is inductively built according to the scheme:

$$
\mathfrak{a} \rightsquigarrow\left(\begin{array}{cc}
\mathfrak{a} & 0  \tag{16}\\
\mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right) \rightsquigarrow\left(\begin{array}{cccc}
\mathfrak{a} & 0 & 0 & 0 \\
\mathfrak{a}_{1} & \mathfrak{a} & 0 & 0 \\
\mathfrak{a}_{2} & 0 & \mathfrak{a} & 0 \\
\mathfrak{a}_{12} & \mathfrak{a}_{2} & \mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right) \rightsquigarrow \ldots,
$$

with repeated, as shown above, $n$-dimensional blocks

$$
\mathfrak{a}=\left(a_{j}^{i}\right), \mathfrak{a}_{1}=\left(a_{j}^{i}\right)_{1}, \quad \mathfrak{a}_{2}=\left(a_{j}^{i}\right)_{2}, \mathfrak{a}_{12}=\left(a_{j}^{i}\right)_{12}, \ldots
$$

Therefore, there follows the general rule: the Jacobian matrix of the transformation of coordinates on the neighborhood $T^{k} U$ is of the form

$$
\left(\begin{array}{cc}
\mathcal{A} & 0  \tag{17}\\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right),
$$

where the block $\mathcal{A}$ is the Jacobian matrix on $T^{k-1} U$ and $\mathcal{A}_{k} \doteq d_{k} \mathcal{A}, k=1,2, \ldots$
In other words, the Jacobian matrix on the neighborhood $T^{k} U$ consists of four blocks, where the Jacobian matrix $\mathcal{A}$ of the neighborhood $T^{k-1} U$ is repeated on the diagonal, the upper-right block is zero, and the left-lower block is the differential of the block $\mathcal{A}$ taking into consideration the $k$-th level, i.e., $\mathcal{A}_{k} \doteq d_{k} \mathcal{A}$.

Formula (17) defines the sequence of matrices (16).
Exercise 16: inversion rule. Show that the inversion of the matrix (17) takes place according to the scheme:

$$
\begin{aligned}
& \left(\begin{array}{cc}
\mathcal{A} & 0 \\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right)=\left(\begin{array}{cc}
\mathcal{E} & 0 \\
\mathcal{A}_{k} \mathcal{A}^{-1} & \mathcal{E}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathcal{A} & 0 \\
0 & \mathcal{A}
\end{array}\right) \rightsquigarrow \\
\rightsquigarrow \quad & \left(\begin{array}{cc}
\mathcal{A} & 0 \\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right)^{-1}=\left(\begin{array}{cc}
\mathcal{A}^{-1} & 0 \\
0 & \mathcal{A}^{-1}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathcal{E} & 0 \\
-\mathcal{A}_{k} \mathcal{A}^{-1} & \mathcal{E}
\end{array}\right),
\end{aligned}
$$

where $\mathcal{E}$ is the identity block. See (4) and

$$
a_{1}=\left(a_{1} a^{-1}\right) a \leadsto a_{1}^{-1}=-a^{-1}\left(a_{1} a^{-1}\right) .
$$

The matrix (17) depends on the point $u_{(k)} \in T^{k} U$. If this point is fixed, then a numeric matrix is defined, but still having the freedom to choose the function $a^{i}$ (or the corresponding jet of the transformation).

Exercise 17: gauge group. Show that all matrices of the form (17), with the point $u_{(k)} \in T^{k} U$ fixed, determine a subgroup of the linear group of order $2^{k} n$,

$$
\mathcal{G}_{k} \subset G L\left(2^{k} n, \mathbb{R}\right)
$$

Prove the existence of the groups $\mathcal{G}_{1}, \mathcal{G}_{2}, \mathcal{G}_{3}$ and extend to $\mathcal{G}_{k}$.
We call the group $\mathcal{G}_{k}$ of matrices (17) with fixed point $u_{(k)} \in T^{k} U$ the gauge group of order $k$ on the manifold $M$. By setting $k=0,1,2,3, \ldots$, we obtain an infinite sequence of gauge groups

$$
\begin{equation*}
\mathcal{G} \rightsquigarrow \mathcal{G}_{1} \rightsquigarrow \mathcal{G}_{2} \rightsquigarrow \mathcal{G}_{3} \rightsquigarrow \ldots \tag{18}
\end{equation*}
$$

Theorem 2. The gauge group of order $k$ is isomorphic to the $k$-th tangent group of the linear group $G L(n, \mathbb{R})$, which, in its turn, is embedded in the linear group $G L\left(2^{k} n, \mathbb{R}\right)$ :

$$
\begin{equation*}
\mathcal{G}_{k} \approx T^{k}(G L(n, \mathbb{R})) \subset G L\left(2^{k} n, \mathbb{R}\right), \quad k=0,1,2, \ldots \tag{19}
\end{equation*}
$$

In this case

$$
\operatorname{dim} G L\left(2^{k} n, \mathbb{R}\right)=\left(2^{k} n\right)^{2} \quad \text { and } \quad \operatorname{dim} \mathcal{G}_{k}=2^{k} n^{2}
$$

Proof. We fix the element $u_{(k)} \in T^{k} M$ of the $k$-th level. Matrices (17) generate a subgroup $\mathcal{G}_{k}$ of the linear group $G L\left(2^{k} n, \mathbb{R}\right)$ (see Exercise 17). The fixing of the point $u_{(k)}$ does not limit the freedom of choice for the element (17) in the group $\mathcal{G}_{k}$. Hence, the group $\mathcal{G}_{k}$ is uniquely defined regardless of the point $u_{(k)} \in T^{k} M$. On the other side, the tangent group $T^{k}(G L(n, \mathbb{R}))$ coincides up to an isomorphism, with the matrix group (17), or $\mathcal{G}_{k}$. This follows from the formulas (3)-(7) and Exercises 12-17, if we assume $G=G L(n, \mathbb{R})$.

Further, in the matrix (17), besides the point $u_{(k)} \in T^{k} U$, there exists the $k$-jet of coordinate transformations $\left(a_{i}^{j}, a_{i_{1} i_{2}}^{j}, \ldots, a_{i_{1} i_{2} \ldots i_{k}}^{j}\right)$. We shall denote as $J_{k}$ the group of such jets at the point $u$ (see Exercise 5). The homomorphism is defined:

$$
\begin{equation*}
\chi_{k}: J_{k} \rightarrow \mathcal{G}_{k} \tag{20}
\end{equation*}
$$

Exercise 18: jets and gauge group. Show that for $k=2$, the mapping $\chi_{2}$ is homomorphic, i.e., to a composition of 2 -jets $\left(a_{k}^{i}, a_{k l}^{j}\right)$ and $\left(b_{j}^{k}, a_{j l}^{k}\right)$ there corresponds the product of matrices $\mathcal{A}_{2}$,

$$
\left(\begin{array}{cc}
a_{k}^{i} & 0 \\
a_{k l}^{i} u_{1}^{l} & a_{k}^{i}
\end{array}\right) \cdot\left(\begin{array}{cc}
b_{j}^{k} & 0 \\
b_{j l}^{k} u_{1}^{l} & b_{j}^{k}
\end{array}\right)=\left(\begin{array}{cc}
a_{k}^{i} b_{j}^{k} & 0 \\
\left(a_{k}^{i} b_{j}^{k}\right)_{l} u_{1}^{l} & a_{k}^{i} b_{j}^{k}
\end{array}\right)
$$

and to the inverse 2 -jet $\left(a_{j}^{i}, a_{j l}^{i}\right)^{-1} \doteq\left(\bar{a}_{j}^{i},-\bar{a}_{s}^{i} a_{k l}^{s} \bar{a}_{j}^{k}\right)$, there corresponds the inverse matrix $\mathcal{A}_{2}^{-1}$

$$
\left(\begin{array}{cc}
a_{k}^{i} & 0 \\
a_{k l}^{i} u_{1}^{l} & a_{k}^{i}
\end{array}\right)^{-1}=\left(\begin{array}{cc}
\bar{a}_{j}^{i} & 0 \\
-\bar{a}_{s}^{i} a_{k l}^{s} \bar{a}_{j}^{k} u_{1}^{l} & \bar{a}_{j}^{i}
\end{array}\right) .
$$

Generalize this to the general case $k$.
Exercise 19: homogeneity of tangent space. Show that the kernel of the homomorphism $\chi_{k}$ is the stabilizer $H_{u_{(k)}}$ of the element $u_{(k)} \in T^{k} M$ in the group $J_{k}$. The tangent space $T_{u_{(k)}}^{k} M$ is identified with the homogeneous space $J_{k} / H_{u_{(k)}}$.

Let us consider once again the gauge groups of the sequence (18). The first group $\mathcal{G}$ is the linear group $G L(n, \mathbb{R})$,

$$
\mathcal{G}=G L(n, \mathbb{R})
$$

The second group $\mathcal{G}_{1}$ is isomorphic to the tangent group $T(G L(n, \mathbb{R}))$. Its elements are block matrices of the form

$$
\left(\begin{array}{cc}
\mathfrak{a} & 0 \\
\mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right), \quad \text { where } \quad \mathfrak{a} \in G L(n, \mathbb{R}) \quad \text { and } \quad \mathfrak{a}_{1} \in \operatorname{gl}(n, \mathbb{R})
$$

The correspondence $\mathcal{G}_{1} \longleftrightarrow T(G L(n, \mathbb{R}))$ is one-to-one. The product of elements in the group $\mathcal{G}_{1}$,

$$
\left(\begin{array}{cc}
\mathfrak{a} & 0 \\
\mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathfrak{b} & 0 \\
\mathfrak{b}_{1} & \mathfrak{b}
\end{array}\right)=\left(\begin{array}{cc}
\mathfrak{a} \mathfrak{b} & 0 \\
(\mathfrak{a b})_{1} & \mathfrak{a b}
\end{array}\right),
$$

reduces to the Leibniz rule in the tangent group $T(G L(n, \mathbb{R}))$,

$$
(\mathfrak{a b})_{1}=\mathfrak{a}_{1} \mathfrak{b}+\mathfrak{a b}_{1}
$$

and the inversion of elements in $\mathcal{G}_{1}$,

$$
\left(\begin{array}{cc}
\mathfrak{a} & 0 \\
\mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right)^{-1}=\left(\begin{array}{cc}
\mathfrak{a}^{-1} & 0 \\
-\mathfrak{a}^{-1} \mathfrak{a}_{1} \mathfrak{a}^{-1} & \mathfrak{a}^{-1}
\end{array}\right)
$$

reduces to the rule

$$
\mathfrak{a}_{1}^{-1}=-\mathfrak{a}^{-1} \mathfrak{a}_{1} \mathfrak{a}^{-1}
$$

This speaks about an isomorphism between the groups $\mathcal{G}_{1}$ and $T(G L(n, \mathbb{R}))$. An inner authomorphism in $\mathcal{G}_{1}$ is generated as allows:

$$
\left(\begin{array}{ll}
\mathfrak{a} & 0 \\
\mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathfrak{b} & 0 \\
\mathfrak{b}_{1} & \mathfrak{b}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathfrak{a} & 0 \\
\mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right)^{-1}=\left(\begin{array}{cc}
\mathfrak{a} \mathfrak{b a} \\
\left(\mathfrak{a b} \mathfrak{a}^{-1}\right)_{1} & \mathfrak{a b a} \mathfrak{a}^{-1}
\end{array}\right),
$$

with the block $\left(\mathfrak{a b a}^{-1}\right)_{1}=\mathfrak{a b}_{1} \mathfrak{a}^{-1}+\mathfrak{a}_{1} \mathfrak{a}^{-1}\left(\mathfrak{a b a}^{-1}\right)-\left(\mathfrak{a b a}^{-1}\right) \mathfrak{a}_{1} \mathfrak{a}^{-1}$, etc.
The following group $\mathcal{G}_{2}$ is isomorphic to the tangent group $T^{2}(G L(n, \mathbb{R}))$. The stair-like structure appears again:

$$
\left(\begin{array}{cccc}
\mathfrak{a} & 0 & 0 & 0 \\
\mathfrak{a}_{1} & \mathfrak{a} & 0 & 0 \\
\mathfrak{a}_{2} & 0 & \mathfrak{a} & 0 \\
\mathfrak{a}_{12} & \mathfrak{a}_{2} & \mathfrak{a}_{1} & \mathfrak{a}
\end{array}\right) \cdot\left(\begin{array}{cccc}
\mathfrak{b} & 0 & 0 & 0 \\
\mathfrak{b}_{1} & \mathfrak{b} & 0 & 0 \\
\mathfrak{b}_{2} & 0 & \mathfrak{b} & 0 \\
\mathfrak{b}_{12} & \mathfrak{b}_{2} & \mathfrak{b}_{1} & \mathfrak{b}
\end{array}\right)=\left(\begin{array}{cccc}
\mathfrak{a} \mathfrak{b} & 0 & 0 & 0 \\
(\mathfrak{a b})_{1} & \mathfrak{a b} & 0 & 0 \\
(\mathfrak{a b})_{2} & 0 & \mathfrak{a b} & 0 \\
(\mathfrak{a b})_{12} & (\mathfrak{a b})_{2} & (\mathfrak{a b})_{1} & \mathfrak{a b}
\end{array}\right),
$$

where

$$
\begin{aligned}
(\mathfrak{a b})_{1} & =\mathfrak{a}_{1} \mathfrak{b}+\mathfrak{a} \mathfrak{b}_{1}, \\
(\mathfrak{a b})_{2} & =\mathfrak{a}_{2} \mathfrak{b}+\mathfrak{a b _ { 2 }}, \\
(\mathfrak{a b})_{12} & =\mathfrak{a}_{12} \mathfrak{b}+\mathfrak{a}_{2} \mathfrak{b}_{1}+\mathfrak{a}_{1} \mathfrak{b}_{2}+\mathfrak{a} \mathfrak{b}_{12}
\end{aligned}
$$

Exercise 20: logarithmic rule for gauge group. Show that while forming the blocks

$$
\mathfrak{a} \rightsquigarrow \mathfrak{a}^{-1} \mathfrak{a}_{1} \rightsquigarrow\left(\mathfrak{a}^{-1} \mathfrak{a}_{1}\right)_{2}=\mathfrak{a}^{-1} \mathfrak{a}_{12}-\mathfrak{a}^{-1} \mathfrak{a}_{2} \mathfrak{a}^{-1} \mathfrak{a}_{1} \rightsquigarrow \ldots
$$

there appears the following property of the logarithmic function:

$$
\ln u \rightsquigarrow \frac{u^{\prime}}{u} \rightsquigarrow \frac{u^{\prime \prime}}{u}-\frac{\left(u^{\prime}\right)^{2}}{u^{2}} \rightsquigarrow \ldots
$$

We shall further denote the Lie algebra of the group $\mathcal{G}_{k}$ by $\overline{\mathcal{G}}_{k}$.

The general scheme is the following. An element of the group $\mathcal{G}_{k}$ is generated according to the principle:

$$
\left(\begin{array}{cc}
\mathcal{A} & 0 \\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right), \quad \text { where } \quad \mathcal{A} \in \mathcal{G}_{k-1}, \quad \mathcal{A}_{k} \in \overline{\mathcal{G}}_{k-1}
$$

The product and the inversion of elements,

$$
\begin{gathered}
\left(\begin{array}{cc}
\mathcal{A} & 0 \\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathcal{B} & 0 \\
\mathcal{B}_{k} & \mathcal{B}
\end{array}\right)=\left(\begin{array}{cc}
\mathcal{A B} & 0 \\
(\mathcal{A B})_{k} & \mathcal{A B}
\end{array}\right), \\
\left(\begin{array}{cc}
\mathcal{A} & 0 \\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right)^{-1}=\left(\begin{array}{cc}
\mathcal{A}^{-1} & 0 \\
\mathcal{A}_{k}^{-1} & \mathcal{A}^{-1}
\end{array}\right),
\end{gathered}
$$

reduce to the rules:

$$
(\mathcal{A B})_{k}=\mathcal{A}_{k} \mathcal{B}+\mathcal{A B} \mathcal{B}_{k}, \quad \mathcal{A}_{k}^{-1}=-\mathcal{A}^{-1} \mathcal{A}_{k} \mathcal{A}^{-1}
$$

The Lie algebra $\overline{\mathcal{G}}_{k-1}$ is identified with the additive subgroup of the matrix group $\mathcal{G}_{k}$, whose matrices have the form:

$$
\left(\begin{array}{cc}
\mathcal{E} & 0  \tag{21}\\
\mathcal{A}_{k} & \mathcal{E}
\end{array}\right)
$$

where $\mathcal{E}$ is the unit block, i.e., the unity of the group $\mathcal{G}_{k-1}$. The product and the inversion of such matrices are performed in the following way:

$$
\begin{gathered}
\left(\begin{array}{cc}
\mathcal{E} & 0 \\
\mathcal{A}_{k} & \mathcal{E}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathcal{E} & 0 \\
\mathcal{B}_{k} & \mathcal{E}
\end{array}\right)=\left(\begin{array}{cc}
\mathcal{E} & 0 \\
\mathcal{A}_{k}+\mathcal{B}_{k} & \mathcal{E}
\end{array}\right) \\
\left(\begin{array}{cc}
\mathcal{E} & 0 \\
\mathcal{A}_{k} & \mathcal{E}
\end{array}\right)^{-1}=\left(\begin{array}{cc}
\mathcal{E} & 0 \\
-\mathcal{A}_{k} & \mathcal{E}
\end{array}\right)
\end{gathered}
$$

All these matrices generate within the group $\mathcal{G}_{k}$ a normal divisor,

$$
\left(\begin{array}{cc}
\mathcal{A} & 0 \\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathcal{E} & 0 \\
\mathcal{B}_{k} & \mathcal{E}
\end{array}\right) \cdot\left(\begin{array}{cc}
\mathcal{A} & 0 \\
\mathcal{A}_{k} & \mathcal{A}
\end{array}\right)^{-1}=\left(\begin{array}{cc}
\mathcal{E} & 0 \\
\mathcal{A} \mathcal{B}_{k} \mathcal{A}^{-1} & \mathcal{E}
\end{array}\right)
$$

An inner automorphism of the group $\mathcal{G}_{k}$ leads to the transformation of the block

$$
\mathcal{B}_{k} \rightsquigarrow \quad \widetilde{\mathcal{B}}_{k}=\mathcal{A B}_{k} \mathcal{A}^{-1}
$$

Under such a transformation, the spectrum of the matrix $\mathcal{B}_{k}$ is preserved. The invariants will be the eigenvalues of this matrix and the corresponding symmetric polynomials, which are coefficients in the Hamilton-Cayley formula.

Exercise 21: Lie algebra of the Lie group. Show that the Lie algebra of an arbitrary Lie group $G$ may be regarded as an additive subgroup and a normal divisor of the tangent group $T G$. Describe the cosets of this normal divisor and the corresponding quotient group of the group $T G$.

Exercise 22: structure constants iterated. The structure constants of a Lie group $G$ have three indices and can be placed into a spacial matrix $a$ ). Prove that the structure constants of the tangent groups $T G, T^{2} G$ and $T^{3} G$ can be similarly put into a spacial matrices of type $b$ ), $c$ ) and $d$ ), respectively.

## 2. Tangent bundles and osculators

2.1. Levels and sector-forms. The tangent functor $T$ iterated $k$ times associates to a smooth manifold $M$ its $k$-fold tangent bundle $T^{k} M$ (the $k$-th level of $M$ ) and associates to a smooth map $\varphi: M_{1} \rightarrow M_{2}$ the graded morphism $T^{k} \varphi: T^{k} M_{1} \rightarrow T^{k} M_{2}$, the $k$-th derivative of $\varphi$. The level $T^{k} M$ has a multiple vector bundle structure with $k$ projections onto $T^{k-1} M$ :

$$
\rho_{s} \doteq T^{k-s} \pi_{s}: T^{k} M \rightarrow T^{k-1} M, \quad s=1,2, \ldots, k
$$

where $\pi_{s}$ is the natural projection $T^{s} M \rightarrow T^{s-1} M$.
Local coordinates in neighborhoods

$$
T^{s} U \subset T^{s} M, s=1,2, \ldots, k, \quad \text { where } \quad T^{s-1} U=\pi_{s}\left(T^{s} U\right)
$$

are determined automatically by those in the neighborhood $U \subset M$, the quantities $\left(u^{i}\right)$ being regarded either as coordinate functions on $U$ or as the coordinate components of the point $u \in U$ :

$$
\begin{aligned}
& U: \quad\left(u^{i}\right), i=1,2, \ldots, n=\operatorname{dim} M \\
& T U: \quad\left(u^{i}, u_{1}^{i}\right), \quad \text { with } u^{i} \doteq u^{i} \circ \pi_{1}, u_{1}^{i} \doteq d u^{i}, \\
& T^{2} U: \quad\left(u^{i}, u_{1}^{i}, u_{2}^{i}, u_{12}^{i}\right),
\end{aligned}
$$

with $u^{i} \doteq u^{i} \circ \pi_{1} \pi_{2}, u_{1}^{i} \doteq d u^{i} \circ \pi_{2}, \quad u_{2}^{i} \doteq d\left(u^{i} \circ \pi_{1}\right), \quad u_{12}^{i} \doteq d\left(d u^{i}\right)$, etc.
We set up the following convention: to introduce coordinates on $T^{k} U$, we take the coordinates on $T^{k-1} U$ and repeat them with an additional index $k$, so that a tangent vector is preceded by its point of origin. This indexing is convenient since at present the symbols with index $s$ become fiber coordinates for the projection $\rho_{s}, s=1,2, \ldots, k$.

Thus, for example, under the projections $\rho_{s}: T^{3} U \rightarrow T^{2} U, s=1,2,3$, the coordinates with indices 1,2 and 3 are each suppressed in turn:

$$
\left.\begin{array}{c}
\left(u^{i} u_{1}^{i} u_{2}^{i} u_{12}^{i} u_{3}^{i} u_{13}^{i} u_{23}^{i} u_{123}^{i}\right) \\
\rho_{1} \swarrow \\
\left(u^{i} u_{2}^{i} u_{3}^{i} u_{23}^{i}\right) \\
\left(u^{i} u_{1}^{i} u_{3}^{i} u_{13}^{i}\right)
\end{array}{\left(4 \rho_{3}\right.}_{i} u_{1}^{i} u_{2}^{i} u_{12}^{i}\right) .
$$

The level $T^{k} M$ is a smooth manifold of dimension $2^{k} n$ and admits an important subspace of dimension $(k+1) n$ called the osculating bundle of $M$ (briefly - osculator) of order $k-1$ and denoted by $\operatorname{Osc}^{k-1} M$. The bundle $\operatorname{Osc}^{k-1} M$ is determined by the equality of the projections

$$
\rho_{1}=\rho_{2}=\ldots=\rho_{k}
$$

meaning that an element of $T^{k} M$ belongs to the bundle Osc $^{k-1} M$ precisely when all its $k$ projections into $T^{k-1} M$ coincide. In this case all coordinates with the same number of lower indices coincide. For example, the first bundle Osc $M$ is determined in $T^{2} U \subset T^{2} M$ by the equation $u_{1}^{i}=u_{2}^{i}$, and the second bundle $\operatorname{Osc}^{2} M$ is determined in $T^{3} U \subset T^{3} M$ by $u_{1}^{i}=u_{2}^{i}=u_{3}^{i}, u_{12}^{i}=u_{13}^{i}=u_{23}^{i}$, etc. The coordinates in $\mathrm{Osc}^{k-1} M$ will be denoted by the derivatives of the coordinate functions on $U$, that is $\left(u^{i}, d u^{i}, d^{2} u^{i}, \ldots, d^{k} u^{i}\right)$.

The immersion $\zeta: \operatorname{Osc} M \hookrightarrow T^{2} M$ and its derivative $T \zeta$ are determined in coordinates by matrix formulas:

$$
\left(\begin{array}{c}
u^{i} \\
u_{1}^{i} \\
u_{2}^{i} \\
u_{12}^{i}
\end{array}\right) \circ \zeta=\left(\begin{array}{c}
u^{i} \\
d u^{i} \\
d u^{i} \\
d^{2} u^{i}
\end{array}\right), \quad\left(\begin{array}{c}
u_{3}^{i} \\
u_{13}^{i} \\
u_{23}^{i} \\
u_{123}^{i}
\end{array}\right) \circ T \zeta=\left(\begin{array}{c}
d u^{i} \\
d^{2} u^{i} \\
d^{2} u^{i} \\
d^{3} u^{i}
\end{array}\right),
$$

$$
T \zeta\left(\frac{\partial}{\partial u^{i}}, \frac{\partial}{\partial\left(d u^{i}\right)}, \frac{\partial}{\partial\left(d^{2} u^{i}\right)}\right)=\left(\frac{\partial}{\partial u^{i}}, \frac{\partial}{\partial u_{1}^{i}}+\frac{\partial}{\partial u_{2}^{i}}, \frac{\partial}{\partial u_{12}^{i}}\right)
$$

The fibres of the bundle Osc $M$ are the integral manifolds of the distribution

$$
\left\langle\partial_{i}^{1}+\partial_{i}^{2}, \partial_{i}^{12}\right\rangle, \quad \text { with } \quad \partial_{i}^{1}+\partial_{i}^{2} \doteq \frac{\partial}{\partial u_{1}^{i}}+\frac{\partial}{\partial u_{2}^{i}}, \quad \partial_{i}^{12} \doteq \frac{\partial}{\partial u_{12}^{i}} .
$$

The functions $\left(u_{1}^{i}-u_{2}^{i}\right)$ vanish on Osc $M$.
Historically, osculating bundles were introduced under various names long before the bundles $T^{k} M$. The systematic study which was initiated 60 years ago by works of V. Vagner [2] has been culminated in recent times in Miron - Atanasiu theory [3]. Meanwhile, the theme of levels $T^{k} M$ remained unjustly neglected for the obvious reason that the multiple fibre bundle structure demands a whole new understanding and new approach (see [1, 4-6]). Attempts such as [7] and the so-called synthetic formulation of $T^{k} M$ [8] made progress in that direction.

While an infinitesimal displacement of the point $u \in M$ is determined by a tangent vector $u_{1}$ to $M$, an infinitesimal displacement of the element $\left(u, u_{1}\right) \in T M$ is determined by the quantities $\left(u_{2}, u_{12}\right)$, representing a tangent vector to $T M$, etc. This interpretation of the elements of $T^{k} M$ allows us to develop the theory of higher order motion. Clearly, the future belongs to these bundles.

White considers on the level $T^{k} M$ or on a $k$-multiple vector bundle certain sectorforms which are functions simultaneously linear on the fibres of all $k$ projections (see [7]). In particular, the sector-forms on $T^{2} U$ and $T^{3} U$ can be written as

$$
\begin{aligned}
& \Phi=\varphi_{i j} u_{1}^{i} u_{2}^{j}+\varphi_{i} u_{12}^{i}, \\
& \Psi=\psi_{i j k} u_{1}^{i} u_{2}^{j} u_{3}^{k}+\psi_{i j}^{1} u_{1}^{i} u_{23}^{j}+\psi_{i j}^{2} u_{2}^{i} u_{13}^{j}+\psi_{i j}^{3} u_{3}^{i} u_{12}^{j}+\psi_{i} u_{123}^{i},
\end{aligned}
$$

with coefficients in $U$. For example, in each term of $\Psi$ the index 1 (or 2 or 3 respectively) appears exactly once. This means that the function $\Psi$ is linear on the fibres of $\rho_{1}\left(\right.$ and $\rho_{2}$ and $\left.\rho_{3}\right)$.

Any scalar function can be lifted from the level $T^{k-1} M$ to the level $T^{k} M$ by $k$ different projections $\rho_{s}: T^{k} M \rightarrow T^{k-1} M$. For example, for the sector-form $\Phi$ (see above) there are three possibilities of lifting to $T^{3} M$ :

$$
\Phi \circ \rho_{1}=\varphi_{i j} u_{2}^{i} u_{3}^{j}+\varphi_{i} u_{23}^{i}, \quad \Phi \circ \rho_{2}=\varphi_{i j} u_{1}^{i} u_{3}^{j}+\varphi_{i} u_{13}^{i}, \quad \Phi \circ \rho_{3}=\varphi_{i j} u_{1}^{i} u_{2}^{j}+\varphi_{i} u_{12}^{i}
$$

Proposition 1. Every exterior $k$-form can be regarded as a sector-form in the sense of White, a scalar function on $T^{k} M$ that is constant on the fibres of $\mathrm{Osc}^{k-1} M$.

Proof. The sector-form $\Phi$ is constant on Osc $M$ if and only if its derivatives vanish on Osc $M$. Thus

$$
\begin{aligned}
\Phi=\varphi_{i j} u_{1}^{i} u_{2}^{j}+\varphi_{i} u_{12}^{i} & \Rightarrow \\
\left(\partial_{i}^{1}+\partial_{i}^{2}\right) \Phi & =\varphi_{i j} u_{2}^{j}+\varphi_{j i} u_{1}^{j}=\left(\varphi_{i j}+\varphi_{j i}\right) u_{1}^{j}-\varphi_{i j}\left(u_{1}^{j}-u_{2}^{j}\right), \\
\partial_{i}^{12} \Phi=\varphi_{i} & \Rightarrow \quad \varphi_{(i j)}=0, \quad \varphi_{i}=0 .
\end{aligned}
$$

If $\Phi$ is an antisymmetric bilinear form then it can be expressed in the coordinates $\left(u^{i}, d u^{i}\right)$ as a 2 -form $\Phi=\varphi_{[i j]} d u^{i} \wedge d u^{j}$. Thus the sector-form $\Phi$ is constant on Osc $M$ if and only if it is a Cartan 2 -form.

If $k=3$ the fibres $\operatorname{Osc}^{2} M$ of dimension $3 n$ are the integral manifolds of the distribution

$$
\left\langle\partial_{i}^{1}+\partial_{i}^{2}+\partial_{i}^{3}, \partial_{i}^{23}+\partial_{i}^{13}+\partial_{i}^{12}, \partial_{i}^{123}\right\rangle .
$$

For the sector-form $\Psi$ (see above) we have

$$
\begin{gathered}
\Psi=\psi_{i j k} u_{1}^{i} u_{2}^{j} u_{3}^{k}+\psi_{i j}^{1} u_{1}^{i} u_{23}^{j}+\psi_{i j}^{2} u_{2}^{i} u_{13}^{j}+\psi_{i j}^{3} u_{3}^{i} u_{12}^{j}+\psi_{i} u_{123}^{i} \Rightarrow \\
\left(\partial_{i}^{1}+\partial_{i}^{2}+\partial_{i}^{3}\right) \Psi=\psi_{i j k} u_{2}^{j} u_{3}^{k}+\psi_{j i k} u_{1}^{j} u_{3}^{k}+\psi_{j k i} u_{1}^{j} u_{2}^{k}+\psi_{i j}^{1} u_{23}^{j}+\psi_{i j}^{2} u_{13}^{j}+\psi_{i j}^{3} u_{12}^{j}, \\
\left(\partial_{i}^{23}+\partial_{i}^{13}+\partial_{i}^{12}\right) \Psi=\psi_{j i}^{1} u_{1}^{j}+\psi_{j i}^{2} u_{2}^{j}+\psi_{j i}^{3} u_{3}^{j}, \\
\partial_{i}^{123} \Psi=\psi_{i} .
\end{gathered}
$$

The derivatives vanish on the fibres $\mathrm{Osc}^{2} M$ when the following conditions hold:

$$
\varphi_{(i j k)}=0, \quad \psi_{i j}^{1}+\psi_{i j}^{2}+\psi_{i j}^{3}=0, \quad \psi_{i}=0
$$

These conditions are necessary and sufficient for the sector-form $\Psi$ to be constant on $\operatorname{Osc}^{2} M$, but not for $\Psi$ to be a Cartan 3 -form. However, every 3 -form $\widetilde{\Psi}=$ $=\varphi_{i j k} d u^{i} \wedge d u^{j} \wedge d u^{k}$ can be regarded as a homogeneous sector-form that is constant on $\operatorname{Osc}^{2} M$.

The argument extends likewise to the cases when $k>3$.
White's theory of sector-forms is much more extensive than that of Cartan exterior forms. In particular, exterior differentiation is an operation on the set of sector-forms that are constant on the osculating bundles.
2.2. Gauge groups on osculating spaces. The action of the gauge group $\mathcal{G}_{k}$ on the $k$-th level $T^{k} M$ extends in a natural way to the osculating bundle $\operatorname{Osc}^{k-1} M$. The diagram from below shows how the block-matrix $4 \times 4$ reduces, for $u_{1}=u_{2}$, to a $3 \times 3$ block-matrix:

$$
u_{1}=u_{2} \quad \Rightarrow\left(\begin{array}{cccc}
a & 0 & 0 & 0 \\
a_{1} & a & 0 & 0 \\
a_{2} & 0 & a & 0 \\
a_{12} & a_{2} & a_{1} & a
\end{array}\right) \rightsquigarrow\left(\begin{array}{ccc}
a & 0 & 0 \\
d a & a & 0 \\
d^{2} a & d a & a
\end{array}\right) .
$$

The blocks of the matrix from the right side are generated in the following way:

$$
\begin{aligned}
a \sim a_{j}^{i}, & \left.\begin{array}{r}
a_{1} \sim a_{j k}^{i} u_{1}^{k} \\
a_{2} \sim a_{j k}^{i} u_{2}^{k}
\end{array}\right\} \rightsquigarrow d a \sim d a_{j}^{i}=a_{j k}^{i} d u^{k}, \\
a_{12} \sim a_{j k l}^{i} u_{1}^{k} u_{2}^{l}+a_{j k}^{i} u_{12}^{k} & \rightsquigarrow d^{2} a \sim a_{j k l}^{i} d u^{k} d u^{l}+a_{j k}^{i} d^{2} u^{k} .
\end{aligned}
$$

The action of the gauge group $\mathcal{G}_{2}$ on the level $T^{2} M$ is obviously transported to the subbundle Osc $M \subset T^{2} M$. While one passes from $T^{2} M$ to Osc $M$ by considering

$$
\left(a_{1}=a_{2}, a_{12}\right) \rightsquigarrow\left(d a, d^{2} a\right), \quad\left(\partial^{1}+\partial^{2}, \partial^{12}\right) \rightsquigarrow\left(\frac{\partial}{\partial(d u)}, \frac{\partial}{\partial\left(d^{2} u\right)}\right)
$$

the transformation of the natural basis on $T^{2} M$ is transported to the transformation of the natural basis on $\operatorname{Osc} M$ :

$$
\left(\begin{array}{lll}
\partial & \partial^{1} & \partial^{2} \\
\partial^{12}
\end{array}\right) \cdot\left(\begin{array}{cccc}
a & 0 & 0 & 0 \\
a_{1} & a & 0 & 0 \\
a_{2} & 0 & a & 0 \\
a_{12} & a_{2} & a_{1} & a
\end{array}\right) \rightsquigarrow\left(\frac{\partial}{\partial u} \frac{\partial}{\partial(d u)} \frac{\partial}{\partial\left(d^{2} u\right)}\right) \cdot\left(\begin{array}{ccc}
a & 0 & 0 \\
d a & a & 0 \\
d^{2} a & d a & a
\end{array}\right) .
$$

In the general case, the action of the group $\mathcal{G}_{k}$ on the level $T^{k} M$ extends in a similar way to the subbundle $\mathrm{Osc}^{k-1} M$.

## Резюме

M. Рахула, В. Балан. Касательные расслоения и калибровочные группы.

Дифференциалы $T^{k} a(k \geq 1)$ диффеоморфизма $a$ гладкого многообразия $M$ индуцируют в слоях расслоений $T^{k} M$, то есть в соответствующих касательных пространствах, линейные преобразования, заключающие в себе действие калибровочной группы $\mathcal{G}_{k}$. Это действие естественным образом распространяется на соприкасающиеся подрасслоения $\mathrm{Osc}^{k-1} M \subset T^{k} M$.

Ключевые слова: диффеоморфизм гладкого многообразия, пространство расслоения, действие калибровочной группы.

## References

1. Ehresmann Ch. Catégories doubles et catégories structurées // C. R. Acad. Sci. - Paris, 1958. - V. 256. - P. 1198-1201.
2. Vagner V.V. Theory of differential objects and foundations of differential geometry // Veblen O., Whitehead J.H.C. The Foundations of Differential Geometry. - Moscow: IL, 1949. - P. 135-223. (in Russian)
3. Atanasiu G., Balan V., Brînzei N., Rahula M. Second Order Differential Geometry and Applications: Miron - Atanasiu Theory. - Moscow: Librokom, 2010. - 250 p. (in Russian)
4. Pradines J. Suites exactes vectorielles doubles et connexions // C. R. Acad. Sci. - Paris, 1974. - V. 278. - P. 1587-1590.
5. Atanasiu G., Balan V., Brînzei N., Rahula M. Differential Geometric Structures: Tangent Bundles, Connections in Bundles, Exponential Law in the Jet Space. - Moscow: Librokom, 2010. - 320 p. (in Russian)
6. Rahula M. Tangent structures and analytical mechanics // Balkan J. Geom. Appl. 2011. - V. 16, No 1. - P. 122-127.
7. White E.J. The Method of Iterated Tangents with Applications in Local Riemannian Geometry. - Boston, Mass.; London: Pitman Adv. Publ. Program, 1982. - 252 p.
8. Bertram W. Differential Geometry, Lie Groups and Symmetric Spaces over General Base Fields and Rings // Memoirs of AMS. - 2008. - No 900. - 202 p.

Поступила в редакцию 17.12.10

[^1]
[^0]:    ${ }^{1}$ Since the time of S. Lie and frequently nowadays, group operators have been called infinitesimal transformations or fundamental vector fields of the group.
    ${ }^{2}$ The matrix $\xi$ plays an essential role in the theory of Lie group representations (see, e.g., S. Lie Theorems).

[^1]:    Rahula, Maido - Doctor of Physics and Mathematics, Professor Emeritus, Faculty of Mathematics and Computer Science, University of Tartu, Tartu, Estonia.

    Рахула, Майдо - доктор физико-математических наук, почетный профессор факультета математики и информатики Тартуского университета, г. Тарту, Эстония.

    E-mail: rahula@ut.ee
    Balan, Vladimir - Doctor of Mathematics, Professor, Faculty of Applied Sciences, University Politehnica of Bucharest, Bucharest, Romania.

    Балан, Владимир - доктор математических наук, профессор факультета прикладных наук Бухарестского политехнического университета, г. Бухарест, Румыния.

    E-mail: vladimir.balan@upb.ro

