

## ON THE EQUATIONS OF DYNAMICS OF THE NON-NEWTON LIQUID WHICH MOVES WITH SLIDING ALONG BED

S.L. Tonkonog

1. In the present article we study the equation

$$u_t = (u^2 \sigma^n + \varepsilon \cdot k(u, \sigma))_x, \quad \sigma = uu_x \quad (1.1)$$

by means of the methods of the approximate group analysis.

Equations of the form (1.1) arise, for example, in simulating one-dimensional flow of a film of non-Newton liquid along a horizontal bed with sliding. In this case, the unknown function  $u(t, x)$  is the thickness of the film at a point  $x$  at the moment  $t$ ,  $n$  is the exponent in rheological law, the second addend in the right-hand side of (1.1) describes the sliding effect,  $\varepsilon$  is a parameter which often can be considered small (in sliding problems in dimensionless equation (1.1) we have  $\varepsilon \approx 0.15$ ). Exact symmetries and invariant solution of equation (1.1) for  $\varepsilon = 0$  (i. e., in absence of sliding) were investigated in [1]–[3] for both one-dimensional case and two-dimensional problems with circular symmetry. Methods which we apply in the present article are based, as in [1]–[3], on the classical group analysis (see [4]) because our investigation is restricted to point groups. This is intrinsic for the class of equations under consideration since even for heat conduction equation with a source its Lie–Becklund group (see [5]) is as a rule (with two exceptions) trivial (see [6]). Their group analysis is reduced to classification of point symmetries (the Lie groups) which was done for nonlinear heat conduction equation without source by L.V. Ovsyannikov (see [7]), and in presence of a source — by V.A. Dorodnitsyn (see [8]). The present article can be considered as a continuation of these works of studying approximate symmetries with a small parameter, their general theory was constructed in [9], [10]. Since equation (1.1) models the flow only up to  $O(\varepsilon^2)$ , it is natural to consider its approximate symmetries with the same accuracy. The classification of equations of the form (1.1), which admit a group of approximate symmetries, is given in Sections 2 and 3. In Section 4 we give approximate-invariant solution with the accuracy  $o(\varepsilon^2)$  for equations obtained in Sections 2 and 3. These equations are of the most interest in the dynamics of non-Newton liquids. For the sake of brevity, further we use the following notation

$$t = y_1, \quad x = y_2, \quad u = y_3, \quad \sigma = y_4, \quad u_t = y_5, \quad u_x = y_6, \quad \sigma_t = y_7, \quad \sigma_x = y_8$$

and then rewrite equation (1.1) into the form

$$F_0^1 + \varepsilon F_1^1 = 0, \quad F_0^2 + \varepsilon F_1^2 = 0, \quad (1.2)$$

where

$$F_0^1 = 2y_3y_6y_4^n + ny_3^2y_4^{n-1}y_8 - y_5, \quad F_0^2 = y_3y_6 - y_4, \\ F_1^1 = \frac{\partial k}{\partial y_3}y_6 + \frac{\partial k}{\partial y_4}y_8, \quad F_1^2 \equiv 0.$$

---

Supported by the Russian Foundation for Basic Research (project no. 94-01-01191).

©1997 by Allerton Press, Inc.

Authorization to photocopy individual items for internal or personal use, or the internal or personal use of specific clients, is granted by Allerton Press, Inc. for libraries and other users registered with the Copyright Clearance Center (CCC) Transactional Reporting Service, provided that the base fee of \$ 50.00 per copy is paid directly to CCC, 222 Rosewood Drive, Danvers, MA 01923.