

GROUP ANALYSIS OF DYNAMIC EQUATIONS OF NON-NEWTON FLUID

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1. In this article, with the use of approximate group analysis developed in [1]–[3], we study the dynamic equation of plane surface of non-Newton fluid subordinate to the Reiner–Rivlin rheological law. It is convenient to write this equation in the form (see [4])

$$u_t = (u^2\Phi(\sigma))_x + \varepsilon F(t, x, u, \varepsilon), \quad \sigma = uu_x. \quad (1.1)$$

Here t is the time, x is the longitudinal coordinate; the unknown function $u(t, x)$ determines the plane surface of fluid, and Φ is an arbitrary function determined by the fluid rheology (usually, in applications, Φ is an odd monotone increasing continuous function), εF is the mass balance of fluid (a near-surface source), ε is a small parameter.

If $\varepsilon = 0$, the unperturbed equation

$$u_t = (u^2\Phi(\sigma))_x, \quad \sigma = uu_x, \quad (1.2)$$

admits the group generated by translations along t and x whose infinitesimal operators are $X_1^0 = \partial_t$ and $X_2^0 = \partial_x$, respectively, and the scaling group whose operator is $X_3^0 = t\partial_t + 2x\partial_x + u\partial_u$.

Note that in [5] the approximate symmetries of (1.1) with Φ being a power function were completely studied. This corresponds to the Ostwald–Lie rheological law for non-Newton fluid (a partial case of the general Reiner–Rivlin rheological law). Also in [5] approximately invariant solutions of the obtained approximately invariant equations were constructed. In [6], [7], for a power function Φ , conditions under which the approximate symmetries of (1.1) are exact, were considered in detail and exact invariant solutions were constructed. In this article we widely use the technique developed in [5]–[7].

2. We shall use results of [1]–[3], therefore the following notation is convenient:

$$\begin{array}{cccccccc} t & x & u & \sigma & u_t & u_x & \sigma_t & \sigma_x \\ y_1 & y_2 & y_3 & y_4 & y_5 & y_6 & y_7 & y_8. \end{array}$$

With this notation, (1.1) is written as

$$F_0^1 + \varepsilon F_1^1 = 0, \quad F_0^2 + \varepsilon F_1^2 = 0, \quad (2.1)$$

where

$$F_0^1 = \frac{\partial}{\partial y_2}(y_3^2\Phi(y_4)) - y_5, \quad F_1^1 = F(y_1, y_2, y_3, y_4), \quad F_0^2 = y_4 - y_3y_6, \quad F_1^2 = 0.$$

By [1], an infinitesimal operator of one-parameter group of approximate up to $O(\varepsilon^2)$ symmetries of (1.1) is generated by an operator of one-parameter group of exact symmetries of unperturbed

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