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PLASMOCHEMICAL METHODS OF PRODUCTION  
AND TREATMENT OF MATERIALS

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## The Film Properties Obtained from Aniline by Plasma Deposition at Atmospheric Pressure

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**Abstract**—The properties of polymer films obtained in plasma of barrier discharge at atmospheric pressure from aniline of amorphous form with simultaneous formation of carbon particles in the films were studied. The threshold of discharge energy density was established ( $\sim 30$  mJ/m), beginning from which a continuous film filled with agglomerates of carbon particles can be obtained. The structure of such films corresponds to the form typical of amorphous polymers. Carbon particles form agglomerates, the concentration of which weakly depends on the discharge energy density and constitutes  $\sim (4-7) \times 10^4$  1/cm<sup>2</sup>, which does not allow increasing the conductivity. It was shown that the main contribution to the conductivity of films obtained from aniline and filled with carbon particles is made by the value of their moisture saturation. It was found that the creep of the material decreases with growth of the energy density, and hardness increases. This is associated with the formation and growth of the number of crosslinks in the samples with an increase in the energy input. The absolute values of Martens hardness correspond to the lower boundary for polyaniline, which is due to the small molecular weight of the obtained polymer.

**Keywords:** atmospheric pressure barrier discharge, carbon particles, conductivity

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### INTRODUCTION

Modern methods of forming polyaniline (PANI) films are directed to increasing their conductivity with simultaneous provision of high strength characteristics such as abrasion resistance and tensile strength [1, 2], which is due to wide prospects of practical application of such films as antistatic transparent coatings of organic glasses (polycarbonate, polymethylmethacrylate), elements of organic electronics for solar batteries and light emitting devices [3, 4], gas identification sensors [5, 6], and humidity sensors [7, 8].

Three main approaches to increasing of the conductivity of PANI films are distinguished: increase in the internal conductivity realized because of self-conductivity of polymer molecules; increase in external conductivity owing to introduction of a conductive

filler, and the combination of these two approaches [9, 10]. It was shown in [10] that protonation of PANI allows stabilizing the positive charge appearing in molecules during oxidation and providing high conductivity of  $\sim 1 - 10$  S/cm. Filling of PANI with carbon particles makes it possible to eliminate the barrier between conducting molecules of the polymer and thereby increase its conductivity. For example, filling of PANI films with carbon nanotubes (MWCNT and SWCNT) during its formation in microwave plasma allows increasing the conductivity from  $10^{-8}$  S/cm for unfilled PANI to  $\sim 1-10$  S/cm when filled with nanotubes by 1.5 wt % [11]. At such a mass concentration, the value of the volume concentration of carbon nanotubes will constitute  $100-1000 \mu\text{m}^{-3}$  under condition of their uniform distribution in the polymer film.

Accounting for the fact that the length of carbon nanotubes reaches  $\sim 1 \mu\text{m}$  and their orientation is equiprobable, at such a concentration, particles will contact each other, forming external conductivity of PANI film.

It is necessary to note that low adhesion of carbon particles to the polymer matrix leads to agglomeration of particles [12], which decreases their volume concentration, and as a result, the PANI conductivity decreases. In addition, in this case, on the boundary particle–polymer, caverns are often formed, which leads to deterioration of the mechanical properties of the polymer film [13]. In [14], an increase in adhesion of carbon nanoparticles to PANI was provided owing to chemical bonds between graphene particles and polyaniline molecules, which allowed providing high stability of the conductivity of protonated PANI films filled with carbon nanoparticles in conditions of an increased humidity ( $\sim 90\%$ ) upon solving the problem of identification of  $\text{CO}_2$  gas. Adhesion of carbon particles of the filler to PANI can also be increased through the formation of PANI film in gas discharge plasma with simultaneous formation of carbon particles in it, as proposed, for example, in [15]. Carbon particles obtained in this way, including nanosized ones, can be charged, which initiates not only the process of polymerization of aniline with the formation of PANI but also the processes similar to protonation of PANI molecules [7].

The aim of this work is to study the structure and properties of the film obtained from aniline in plasma of barrier gas discharge of atmospheric pressure with simultaneous formation of carbon particles in it.

### SAMPLING METHOD

The formation of an aniline film with the simultaneous formation of carbon particles in it was carried out in an AC barrier corona discharge according to the procedure described in [15–17] on a glass slide (SP-7102 glass with thickness of 1 mm). The strength of electric field was regulated in the range  $\sim 6$ – $12 \text{ kV/cm}$  with the accuracy of  $\pm 0.5 \text{ kV/cm}$ . The formation of PANI film occurred from aniline vapor in an argon atmosphere. The pressure in the discharge chamber was  $1040 \pm 5 \text{ GPa}$ . After obtaining the film, it was held in vacuum at a pressure not greater than  $10^{-1} \text{ Torr}$  and temperature of  $95 \pm 5^\circ\text{C}$  for 4 h with a further temperature decrease to  $20$ – $25^\circ\text{C}$  for 20 h. In the experiments, four types of PANI film samples were obtained at the following values of energy density released in the volume of the charge channel (streamer): sample 1,  $15$ – $20 \text{ mJ/m}$ ; sample 2,  $24$ – $27 \text{ mJ/m}$ ; sample 3,  $30 \pm 2 \text{ mJ/m}$ ; sample 4,  $45 \pm 2 \text{ mJ/m}$ .

The surface morphology of the films obtained from aniline in plasma were studied on a Carl Zeiss Merlin scanning electron microscope. A graphite film with a

layer of  $\sim 10 \pm 2 \text{ nm}$  was deposited on the surface of the samples to drain the charge from the samples during their analysis.

The film profile was studied on a Bruker DektakXT™ profilometer (stylus radius of  $12.5 \mu\text{m}$ ) with an error not worse than 10 nm.

Transmission electron microscopy (TEM) was carried out on a Zeiss Libra 120 microscope with a built-in OMEGA filter under the following conditions: accelerating voltage, 120 kV; electron source, LaB6. The analysis of electron diffraction reflections was carried out according to the procedure given in [25–28]. Absorption spectra were obtained on a StellarNet EPP2000 Spectrometer with a resolution of 0.5 nm.

The dielectric properties of the film samples were studied by dielectric spectroscopy at direct current according to GOST R 50499-93 (IEC 93-80) and at alternating current using a Novocontrol BDS-80 dielectric spectrometer. The measurement error of specific conductivity did not exceed 10%. Measurements of specific conductivity on alternating current were carried out at a temperature of  $\sim 23 \pm 2^\circ\text{C}$  and an applied voltage of  $0.5 \pm 0.05 \text{ V}$ .

The Martens microhardness was measured on a Shimadzu DUH-211S ultramicrotester according to ISO 14577-1 using a Berkovich indenter (triangular indenter with apex angle of  $115^\circ$ ), in much the same way as in [38].

## RESULTS AND DISCUSSION

### *Study of the Morphology of the Films by the Method of Electron Microscopy*

The surface of the films obtained from aniline in plasma (Fig. 1) corresponds to the form characteristic of amorphous polymers in contrast to the granular film structure obtained by electrochemical deposition or vacuum spraying [5, 10].

Carbon particles forming in plasma create agglomerates, which are distributed fairly uniformly along the surface of the samples and have an average size of  $4$ – $6 \mu\text{m}$ . No significant dependence of the average size of the agglomerates on the level of energy density was found. The change in their concentration has a weak tendency to increase with growth of energy input to the discharge and is estimated by the value of  $(4$ – $7) \times 10^4 \text{ cm}^{-2}$ . At such a concentration, carbon agglomerates do not come in contact and most likely should not influence the electrical conductivity of the film samples.

The relief profile of the films is shown in Fig. 2.

At energy densities lower than  $\sim 30 \text{ mJ/m}$ , a non-continuous film is formed. With growth of energy density, the rate of film growth increases from  $\sim 1.2 \pm 0.2 \mu\text{m/min}$  (for energy density of  $\sim 25 \text{ mJ/m}$ ) to  $\sim 3 \pm 0.5 \mu\text{m/min}$  (for  $\sim 45 \text{ mJ/m}$ ); disparate regions begin to combine with the formation of a continuous

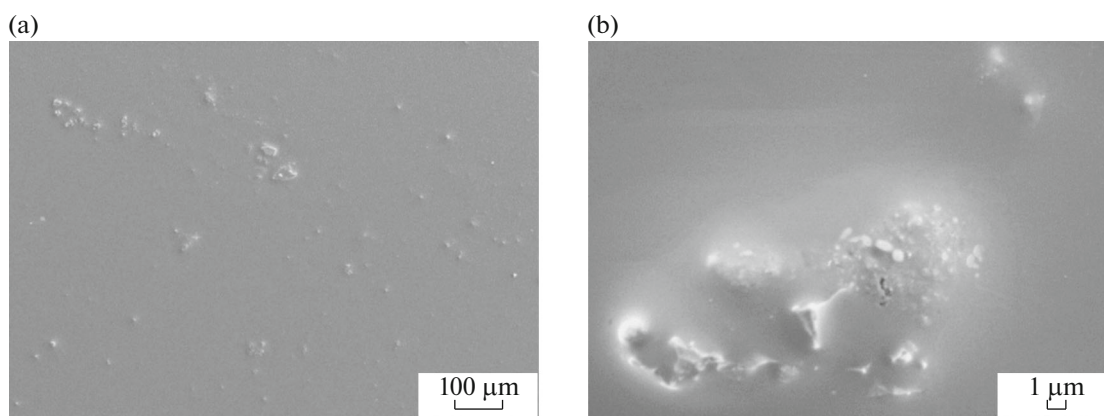


Fig. 1. Typical morphology of the film, sample 4.

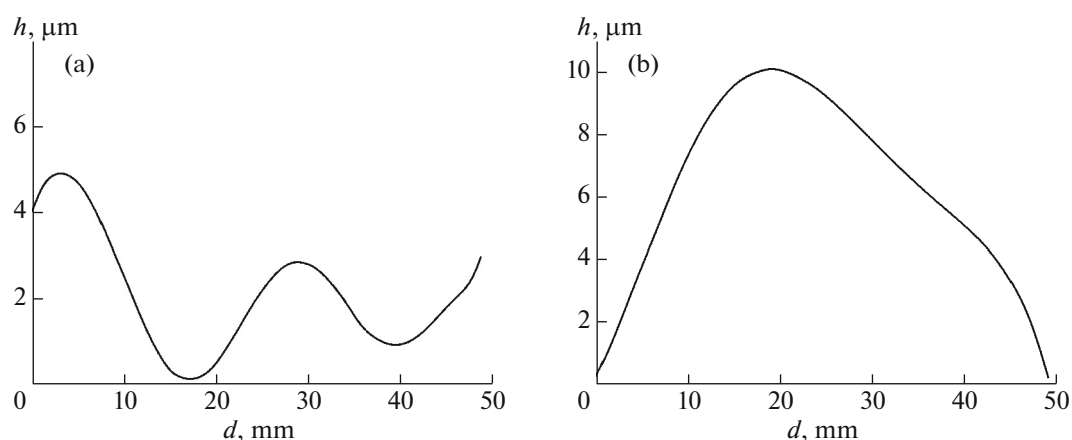


Fig. 2. Typical profiles of the film surface ( $h$ —height,  $d$ —width of films): (a) sample 1, (b) sample 3.

film. Beginning from the energy density of  $\sim 45$  mJ/m, the growth rate of the polymer film constitutes  $\sim 3 \pm 0.5$   $\mu\text{m}/\text{min}$  and weakly depends on further increase in energy density. This can be explained by concurrence of the processes of film formation and its destruction in gas discharge plasma [16] beginning from a specific energy density value in the discharge. It is necessary to note that destruction of aniline during the formation of the film leads to formation of both low molecular organic products and carbon particles, which influence the physicochemical properties of the coating. In the case under consideration, an increase in discharge energy density, apparently, leads to an increase in the fraction of low molecular organic products in the film; an increase in the concentration of carbon particles with growth of energy density is not observed.

In this way, in the range of discharge energy densities of 15–27 mJ/m, a film with the least fraction of low molecular products is formed.

Absorption spectra (Fig. 3) of the film samples have characteristic peaks in the regions of 332, 420, 900, and 970 nm.

The absorption peak at a wavelength of 332 nm (Fig. 3a) corresponds to the benzene ring [14]. The absorption peak at  $\lambda = 420$  nm is pronounced for samples 2–4 and is typical of protonated polyaniline (PANI) [14, 18]. It is known that PANI films have good moisture absorption [19]; therefore, after removing the film samples from the vacuum chamber into the atmosphere, the films were saturated with moisture, as evidenced by a weak peak in the region of 970 nm (Fig. 3b), which is characteristic of hydroxyl groups [20].

In the experiments, the humidity in the room was controlled by a hygrometer with an error of no more than 5%; its value was in the range of 75–85%. The presence of a significant amount of hydroxyl groups in the film contributed to the protonation of PANI films according to the mechanism described in [18]: a weak peak in the region of 900 nm is associated with the  $\pi$ -polaron [14, 21]. It should be noted that there are no peaks at 420, 980, and 900 nm in the transmission spectrum of sample 1. Perhaps this is due to the fact that, at the appropriate value of the energy density, the

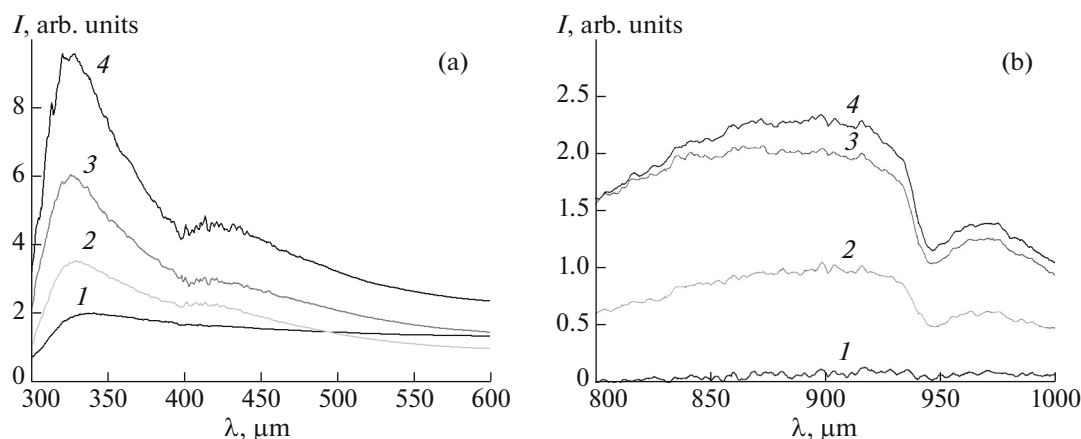


Fig. 3. Absorption spectra of the films: (1–4) samples 1–4, respectively.

formation of radicals necessary for the formation of PANI does not occur.

#### *Identification of Allotropic Form of Carbon Particles by TEM Method*

In a corona discharge, two regions are distinguished in which heat is released owing to the flow of electric current: these are the regions of ionization (corona sheath) and charge drift (streamers) [22, 23]. The plasma temperature in the corona sheath reaches  $\sim 4000^\circ\text{C}$  for the operating modes of the experimental setup used [15, 23]. The average temperature of the medium in the volume of the streamer in the same modes is  $\sim 100\text{--}2200^\circ\text{C}$ . At such temperatures, the destruction of aniline monomer macromolecules [24] and the formation of carbon particles of both amorphous and crystalline allotropic forms are possible.

For investigation of carbon particles by the TEM method (Fig. 4), a suspension was prepared by dissolving a polymer film in acetone (99.9%). After that, the obtained solution was held in an ultrasonic disperser for not less than 10 min (frequency of 44 kHz), the power of ultrasonic impact was  $\sim 10$  W, and the solution was applied in the volume of  $2\ \mu\text{L}$  on a 3 mm copper grid with a Formvar/Carbon substrate (TedPella).

Diffraction analysis of the sample 3 showed a chaotic dislocation of reflections (Fig. 4f), which is due to agglomeration of carbon particles. Nevertheless, reflections corresponding to hexagonal structure in carbon particles can be distinguished in the obtained diffraction patterns, as well as in Fig. 4d:  $d_{100}/d_{110} \approx 1.73$  [28].

The diffraction patterns of electrons (Fig. 4b) are close to those on multilayer carbon nanoparticles, for example, onion-like carbon [26, 28], and in Fig. 4f, on multilayer graphene structures (multilayered graphitic sheets) or soot crystallites [29, 30].

The results of TEM show the formation of both carbon particles with amorphous (Fig. 4b) and crystalline (Figs. 4d, 4f) allotropic forms. In this case, the fraction of particles with crystalline form increases with growth of energy input to the discharge, which is associated with the growth of temperature in the corresponding areas of the corona discharge.

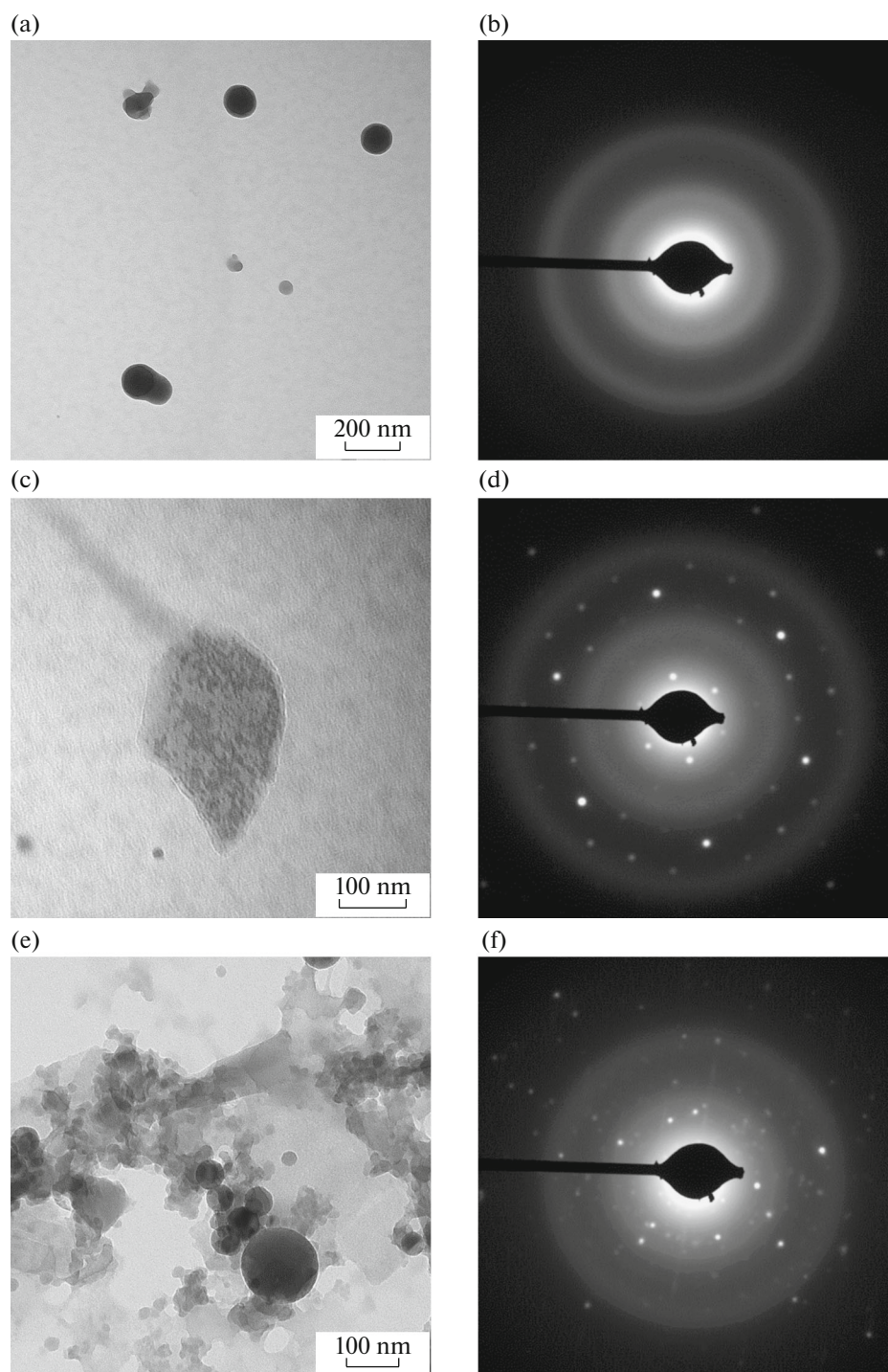
#### *Investigation of Dielectric Properties of the Films by the Method of Dielectric Spectroscopy*

The main contribution to the conductivity of the obtained films from aniline in plasma is made by the value of their moisture saturation (Fig. 5). The volume concentration of carbon particles in the films is insufficient for increasing the external conductivity, and partial protonation with hydroxylic groups [20] does not allow increasing the conductivity. Reduction of the conductivity with an increase in the charge power is apparently associated with a decrease in the limiting moisture saturation of the film by means of formation of crosslinks between molecules and growth of the density of the polymer [31]. After drying of the film samples at pressure of  $\sim 10^{-1}$  Torr and temperature of  $\sim 23 \pm 2^\circ\text{C}$ , the conductivity was no greater than  $10^{-8}$  S/cm.

It is necessary to note that, in a number of works, for example, [14, 32–34], the results of effective use of various sensors based on PANI films filled with carbon nanoparticles at an increased humidity of  $\sim 90\%$  are presented.

#### *Investigation of Mechanical Properties*

The Martens microhardness was determined by the slope of the applied load increase curve and the change in the depth of the indenter impression at constant loading ( $C_{IT}$ , %—the parameter characterizing the creep of the material). Several independent measurements were carried out in different areas of the

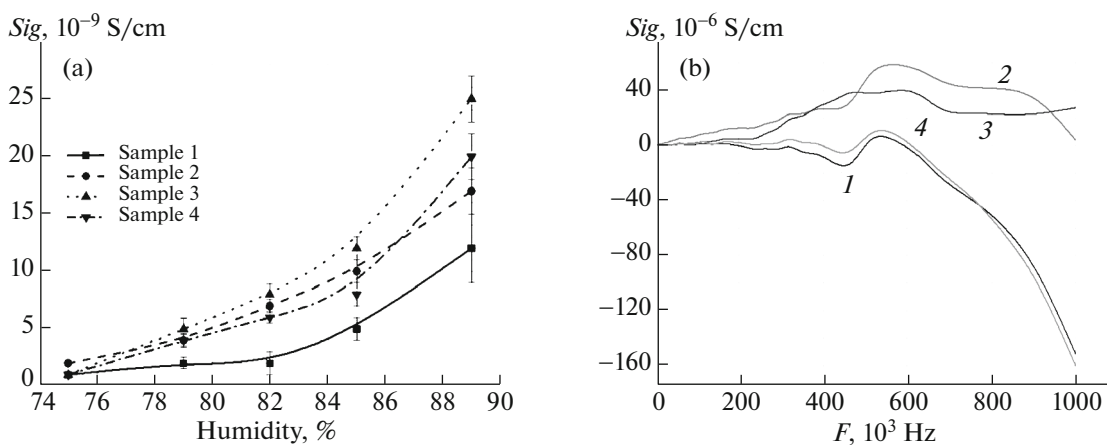


**Fig. 4.** Structures (a, c, e) and electron diffraction (b, d, f) of carbon particles obtained during the formation of PANI film (sample 3).

samples with further averaging of the obtained values (Table 1). The indenter indentation force was 0.17 mN; it was chosen in such a way as to minimize the influence of the substrate on the measurement results [38]; the exposure time at maximum load was

28 s; the parameter  $\eta_{IT}$  characterizes the elasticity of the material [36].

With growth of energy density, the creep of the material of the samples ( $C_{IT}$ ) decreases somewhat, and the hardness increases. This may be caused by the for-



**Fig. 5.** (a) Dynamics of the conductivity of the films at direct current; (b) dielectric spectroscopy of sample 2 at humidity values: (1) 75%, (2) 89%, (3) 85%, (4) 79%.

mation and growth of the number of crosslinks in the samples with an increase in energy input [22, 37]. The absolute values of Martens hardness correspond to the lower boundary for polyaniline [38, 39], which is probably due to the small molecular weight of the obtained polymer.

## CONCLUSIONS

The possibility of obtaining films from aniline of an amorphous form with simultaneous formation of carbon particles in them in the plasma of an AC barrier corona discharge at atmospheric pressure was experimentally confirmed.

The fraction of carbon particles of crystalline alloptropic shape increases with growth of the discharge energy density, which is associated with an increase in temperature in the streamers and the corona discharge sheath. Carbon particles form agglomerates, the concentration of which weakly depends on the discharge energy density and constitutes  $\sim(4-7) \times 10^4 \text{ cm}^{-2}$ , which does not allow increasing the conductivity of the films by means of increasing the external conductivity.

The main contribution to the conductivity of the films filled with carbon particles is made by their moisture saturation, which is determined by the structure of the coating and its continuity. It was established that the surface structure of the obtained films corresponds to the form characteristic of amorphous polymers. Continuity of the coating is achieved at energy densities higher than  $\sim 30 \text{ mJ/m}$ , and the average rate of film growth changes from  $\sim 1.2 \pm 0.2 \text{ }\mu\text{m/min}$  (for energy density of  $\sim 25 \text{ mJ/m}$ ) to  $\sim 3 \pm 0.5 \text{ }\mu\text{m/min}$  (for  $\sim 45 \text{ mJ/m}$ ).

The microhardness of the films increases with growth of energy input to the discharge, which is due to the formation and growth of the number of crosslinks in the samples. The observed peaks in the regions of 420, 900, and 970 nm in the absorption spectra indicate the protonation of polyaniline films, which is initiated by hydroxyl groups according to the mechanism described in [20]. The corresponding peaks are observed only for continuous films, since at an energy density value of less than  $\sim 30 \text{ mJ/m}$ , the formation of radicals necessary for the formation of PANI does not occur.

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**Table 1.** The results of measurements of the Martens microhardness

Sample	Martens hardness $HMs$ , MPa	Material creep $C_{IT}$ , %	Material elasticity $\eta_{IT}$ , %
1	$17.0 \pm 0.5$	$38.6 \pm 0.5$	$2.7 \pm 0.2$
2	$17.5 \pm 0.3$	$37.8 \pm 0.3$	$2.8 \pm 0.4$
3	$18.0 \pm 0.4$	$37.2 \pm 0.3$	$3.2 \pm 0.3$
4	$18.6 \pm 0.5$	$36.8 \pm 0.3$	$3.0 \pm 0.4$

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