

2 **1** Size and Shape of Steady Seawater Intrusion and 3 Sharp-Interface Wedge: The Polubarinova-Kochina 9 **4** **2** Analytical Solution to the Dam Problem Revisited

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6 **Abstract:** Rescaling of the geometrical sizes and the value of hydraulic conductivity in the classical problem of steady two-dimensional (2D)
7 potential seepage through a rectangular earth dam with an empty tailwater is shown to result in a mathematically equivalent problem of
8 seawater intrusion with a sharp interface into a confined horizontal aquifer, which discharges fresh groundwater to the sea through a vertical
9 segment of the beach. The shape of the interface, the vertical and horizontal sizes of the static intrusion wedge, and its cross-sectional area are
10 written in an explicit form, using the Polubarinova-Kochina formulas, rectified. The densities of the two liquids and the aquifers' hydraulic
11 conductivity and thickness, as well as the incident hydraulic gradient serve as input parameters. With reduction of the incident groundwater
12 gradient far upstream from the intrusion zone (due to, e.g., freshwater abstraction by wells), the sizes of the wedge rapidly increase. The
13 analytical solution has been validated with recent sand tank experiments. DOI: [10.1061/\(ASCE\)HE.1943-5584.0001385](https://doi.org/10.1061/(ASCE)HE.1943-5584.0001385). © 2016 American
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15 **Author keywords:** Seawater intrusion; Steady potential phreatic flow; Sharp-interface model; Earth dam problem; Exact solution.

16 Introduction

17 **10** Fresh groundwater discharge as submarine springs or outseeps is
18 important for the global hydrological balance and catchment-scale
19 assessments of travel times of groundwater particles, hydrogeo-
20 chemistry of coastal sea water and discharging aquifers, ecology
21 of coral reefs and fish in Oman affected by groundwater-imported
22 nutrients, paleohydrogeology-anthropology, hydrology of global
23 climate changes, and planning of wellfield operations in coastal
24 zones, among others (e.g., Burnett et al. 2006; Faure et al.
25 2002; Ferguson and Gleeson 2012; Hoefel and Evans 2001; Sherif
26 et al. 2014; Taniguchi et al. 2002; Uchiyama et al. 2000; Zektser
27 and Loaiciga 1993). Seawater intrusion (SWI) in pristine (anthro-
28 pogenically intact) aquifers is conceptualized as a wedge (tongue)
29 of a relatively dense seawater encroaching along the aquifer bottom
30 [e.g., Fig. 1 of Burnett et al. (2006), Fig. 1 of Strack and Ausk
31 (2015), and Fig. 1(a)] against the direction of groundwater dis-
32 charge. SWI, especially with upstream freshwater pumping (the
33 wedge is then bumped in shape and blurred) has a detrimental effect
34 on water supply from coastal aquifers in Oman and other Gulf
35 countries, especially on agricultural and municipal wells. Modeling
36 of SWI is carried out by sharp interface and variable density codes,
37 both analytically and numerically (e.g., Al-Bitar and Ababou 2005;
38 Bakker 2014; Bear and Dagan 1964; Bereslavski 2007; Detournay
39 and Strack 1988; Cheng and Ouazar 1999; De Josselin De Jong and
40 Van Duijn 1984; Glover 1959; Hocking and Forbes 2004; Kacimov
41 and Sherif 2006; Kacimov et al. 2009; Kashef 1983; Kourakos and

Mantoglou 2015; Lu et al. 2015; Llopis-Albert and Pulido-
Velazquez 2015; Mazi et al. 2014; Paster and Dagan 2008; Sherif
et al. 2012; Strack 1989; Strack and Ausk 2015; Werner et al.
2012). Field studies of SWI are based on surface and downhole
geophysics. Laboratory experiments aimed at measuring the sizes
of the wedge [$|B_i D_{1i}|$ and $|A_i D_{1i}|$ in Fig. 1(a)] were carried out in
sand-filled boxes (e.g., Bertorette 2014; Chang and Clement 2012;
Goswami and Clement 2007).

The main question in mathematical models of SWI is what
size of wedge is a quasi-triangle $A_i B_i D_{1i} A_i$ (the subscript i indi-
cates intrusion) in a vertical cross section of Fig. 1(a), i.e., what
are the length and height of the wedge? This paper answers this
question using an exact analytical solution for a sharp-interface
steady-state, Darcian fresh groundwater discharge over a static
saline wedge.

Two-Element Freshwater Discharge into Coastal Aquifer and Seepage through a Rectangular Dam

Similar to Kashef (1983), a confined, isotropic, homogeneous aqui-
fer [Fig. 1(a)] of thickness H_{1i} , with an impermeable caprock and
bedrock, $E_{1i} D_i$ and $E_{2i} D_{1i}$ (two horizontal rays) as their boundaries
is considered. Hydraulic conductivity of the aquifer is k_i . Fresh
groundwater of density r_f moves from the left (Jabal Al-Akdar
mountains in Oman) to the right (shore of the Gulf) and discharges
into the sea as a submarine outlet through an outflow face $D_i A_i$
of the beach. The density of seawater is r_s , $r_s > r_f = 1,000 \text{ kg/m}^3$.
Unlike Kashef (1983) hydrostratigraphy, the aquifer in Fig. 1(a) is
not hydraulically commingled with the superjacent or subjacent
aquifers. A static SWI wedge is bounded from the right by a vertical
segment $A_i D_{1i}$, and from above by a sharp interface $B_i A_i$.

The origin of a Cartesian (x_i, y_i) coordinate system is selected as
point D_i . The tip B_i of the wedge is at $x_i = -l_i$ ($l_i > 0$ is un-
known). The hydraulic head $h_i(x, y)$ (a harmonic function within
the flow domain) of the moving freshwater is counted from point
 D_i , i.e., $h_i(0, 0) = 0$. The tip A_i is at the depth $y_i = -H_{0i}$ ($H_i > 0$ is

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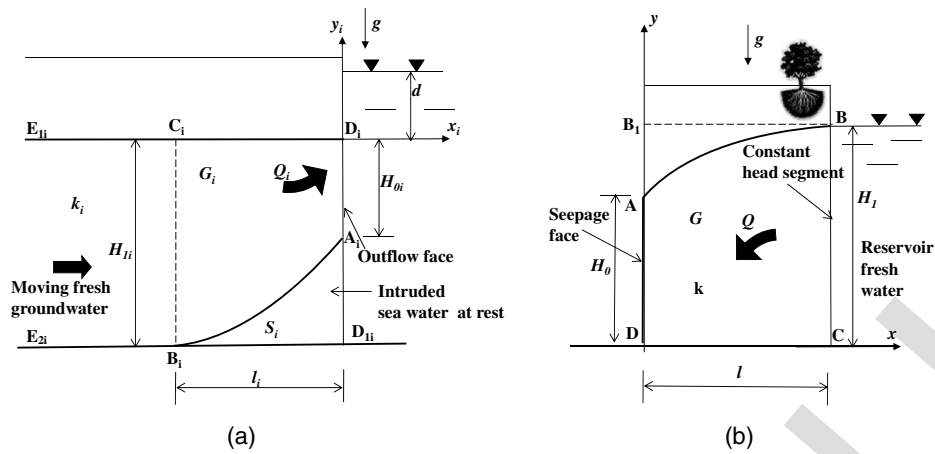


Fig. 1. (a) Vertical cross section of seawater intrusion zone; (b) vertical cross section of seepage in an earth dam

also unknown). The discharge rate of two-dimensional (2D) flow in the plane $x_i D_i y_i$ (per unit length in the direction perpendicular to this plane), $Q_i = \text{const} > 0$ (m^2/s), is given.

On a catchment scale the value of Q_i is often known from the hydrological balance or piezometric data upstream of the SWI zone. In the laboratory experiments of Goswami and Clement (2007) (GC), this value was directly measured.

The goal of this paper is to analyze how the shape of $A_i B_i$ in Fig. 1(a) depends on Q_i , in particular, how l_i , H_{0i} , and the cross-sectional area, S_i , of the wedge in Fig. 1(a) (area between $A_i B_i$ and $B_i D_{1i}$) vary with Q_i .

Kacimov and Obnosov (2001) (KO) gave a full analytical solution to the 2D flow problem in Fig. 1(a), even for a tilted beach face $D_i D_{1i}$. This solution has not, however, won the hearts and minds of groundwater engineers because of its apparent mathematical complexity. A simplified version of this solution is presented here, which combines the rigor of the full 2D potential model with the classical analytical solution of Polubarinova-Kochina (1962, 1977) (PK) to the so-called earth-dam problem shown in Fig. 1(b). KO mentioned the commonality of the problems in Figs. 1(a and b) but—to the best of the authors' knowledge—nobody exploited this analogy in practical groundwater hydrology. Thus, for the sake of methodological lucidity, this paper repeats here what is well-known to geotechnical engineers and applied mathematicians (Crank 1984; Craster 1994; Hornung and Krueger 1985) as the PK dam problem and its solution. This paper illustrates how this solution can be rescaled to the problem of SWI in Fig. 1(a).

A rectangular dam, whose vertical cross section is shown in Fig. 1(b), has the width l and is made of a homogenous isotropic soil of hydraulic conductivity k . The dam stands on an impermeable horizontal foundation CD . The upper pool is filled with a fresh water up to a level H_1 , which is counted from the Dx -axis of a Cartesian coordinate system xy . The vertical face BC is a constant-head boundary. The tailwater is empty and seepage is from the right to the left. The phreatic surface BC is a sharp interface that separates a fully saturated flow domain beneath from an absolutely dry soil above BC . The dry soil triangle BAB_1 in Fig. 1(b) is an analogue of the wedge in Fig. 1(a). Similar to ignoring dispersion and diffusion in the problem of Fig. 1(a), the PK model of seepage in Fig. 1(b) ignores capillarity of the soil. The outlet vertical segment AD in Fig. 1(b) is a seepage face (i.e., an isobar of atmospheric pressure). The shape of AB in Fig. 1(b), in particular the locus of point A (H_0), is a part of the mathematical solution,

as well as the flow rate Q . The hydraulic head, $h(x, y)$, in Fig. 1(b) is counted from point D .

Analytical Solution

As in Kashef (1983), it was assumed that the dashed vertical line $B_i C_i$ in Fig. 1(a) is a line of constant head h_i . This is, of course, only an approximation in terms of the full solution of KO. With this assumption, flow in Fig. 1(a) decouples into two analytic elements: a purely confined trivial one-dimensional (1D) flow in the half-strip $E_{1i} C_i B_i E_{2i}$ and a free-boundary 2D flow on the right of $C_i B_i$. The Dupuit-Forchheimer (DF) approximation (e.g., Bakker 2014; Koussis et al. 2015) adopted by Kashef (1983) was not assumed, and the segment $D_i A_i$ in Fig. 1(a) was not assumed to be a constant hydraulic head boundary. This segment is an outflow boundary as in Strack and Ausk (2015). Unlike Kashef (1983), for the right fragment $B_i C_i D_i A_i B_i$ (a quasi-trapezium G_i) in Fig. 1(a), a potential sharp-interface model of PK was utilized. As usual, a complex potential $w_i = \phi_i + i \psi_i$ is introduced, where i is an imaginary unit and $\phi_i = -k_i h_i$, which is the velocity potential according to the Darcy law $\mathbf{V}_i = -k_i \nabla h_i$, where \mathbf{V}_i is the Darcian velocity vector and ψ_i is a stream function, complexly conjugated with ϕ_i . Parameter ψ_i will be counted from the streamline $C_i D_i$.

Then the boundary value problem (BVP) for flow in Fig. 1(a) is

$$\begin{aligned} B_i C_i: \varphi_i &= -k_i \delta H_{1i}, & \text{at } x_i &= -l_i, & -H_{1i} < y_i < 0 \\ C_i D_i: \psi_i &= 0, & \text{at } y_i &= 0, & -l_i < x_i < 0 \\ D_i A_i: \varphi_i - k_i \delta y_i &= 0, & \text{at } x_i &= 0, & -H_{0i} < y_i < 0 \\ A_i B_i: \psi_i &= -Q_i, & \varphi_i - k_i \delta y_i &= 0 & \text{at } -l_i < x_i < 0, \\ & & -H_{0i} > y_i > -H_{1i} & & \end{aligned} \quad (1)$$

where $\delta = (r_s - r_f)/r_f$. The fresh water head at point B_i [the first line in Eq. (1)] follows from the Pascal law (see KO) and definition of pressure $p_i = r_f g(-\phi_i/k_i - y_i + r_s/r_f d)$, where d is the depth of seawater above point D_i in Fig. 1(a). The boundary condition along $D_i A_i$ is neither constant head nor constant flux, although in regional-scale numerical models like MODFLOW and SEAWAT these simplified conditions are also used (e.g., Motz and Sedighi 2009).

A standard trick (e.g., Kashef 1983) is to rescale the geometric and flow variables of Fig. 1(a) as $(x, y, \phi, \psi, k) = (-\delta x_i, -\delta y_i, \delta \phi_i, \delta \psi_i, \delta k_i)$, where $w = \varphi + i\psi$ is the complex

153 potential of flow in Fig. 1(b). In other words, G_i in Fig. 1(a) is
 154 mapped symmetrically with respect to point D, and stretched to
 155 get G . A porous medium of conductivity k_i is also replaced by
 156 a medium of conductivity k .
 157 Then BVP in Eq. (1) is reduced to

$$\begin{aligned} \text{BC: } \varphi &= -kH_1, & x &= l, & H_1 &> y > 0 \\ \text{CD: } \psi &= 0, & y &= 0, & l &> x > 0 \\ \text{DA: } \varphi + ky &= 0, & x &= 0, & H_1 &> y > 0 \\ \text{AB: } \psi &= -Q, & \varphi + ky &= 0 \end{aligned} \quad (2)$$

158 The rescaled BVP in Eq. (2) exactly corresponds to the dam
 159 flow problem in Fig. 1(b). A full solution to the BVP in Eq. (2)
 160 is given in PK and Crank (1984). Therefore, the back-scaling
 161 immediately solves the BVP in Eq. (1). Rephrasing, if G is flipped
 162 in Fig. 1(b) (including the phreatic surface) over the center of sym-
 163 metry D and stretched 33 times (the density of seawater in the Gulf
 164 corresponds to $\delta \approx 0.03$), the result is G_i in Fig. 1(a) (including the
 165 sharp interface).

166 The flow rates Q and Q_i are calculated by the Charny formula
 167 (PK)

$$Q = k \frac{H_1^2}{2l}, \quad Q_i = \frac{Q}{\delta} = k_i \delta \frac{H_{1i}^2}{2l_i} \quad (3)$$

168 which has been recently extended to layered aquifers by Strack and
 169 Ausk (2015).

170 Eq. (3) for the discharges in both SWI [Fig. 1(a)] and dam
 171 [Fig. 1(b)] problems are exact. Q and Q_i given by Eq. (3) coincide
 172 with those derived from the DF model (see PK for details) but the
 173 sharp interface and phreatic surface clearly do not coincide in the
 174 exact 2D and approximate 1D (DF) models.

175 Eq. (3) is used to calculate the areas of G_i in Fig. 1(a) and G in
 176 Fig. 1(b). Then for the wedge area

$$S_i = H_{1i}l_i + \int_{-l_i}^0 y_{iBA}(x_i)dx_i = \frac{H_1l}{\delta^2} - \frac{1}{\delta^2} \int_0^l y_{AB}(x)dx \quad (4)$$

177 The last integral in Eq. (4) [the saturated area of G in Fig. 1(b)]
 178 requires some effort to evaluate. Unfortunately, in both PK
 179 (1962, 1977) there are numerous typos and an ambiguous statement
 180 on the shape of AB in Fig. 1(b). Namely, after Eq. (10.41) in PK
 181 (1977, Chapter 7) [the same mistake is in PK (1962)], the authors
 182 incorrectly wrote that the parametric equation of the free surface
 183 involves an arbitrary constant. PK suggests equating this
 184 constant to 1. In reality, this constant is not arbitrary but has to
 185 be determined from PK [(1962, 1977, Eqs. (10.34) and (10.35)].
 186 These two equations are rewritten in a dimensionless form as
 187 one equation

$$l^* = \frac{l_i}{H_{1i}} = \frac{l}{H_1} = \frac{\int_0^{\pi/2} \frac{K[\beta \sin^2 \chi]}{\sqrt{1-\beta \sin^2 \chi}} d\chi}{\int_0^{\pi/2} \frac{K[\beta+(1-\beta) \sin^2 \chi]}{\sqrt{\beta+(1-\beta) \sin^2 \chi}} d\chi} \quad (5)$$

188 where K = complete elliptic integral of the first kind; and
 189 $0 \leq \beta \leq 1$ = parameter [the affix of a conformal mapping whose
 190 preimage is point D in Fig. 1(b)], to be determined. The notations
 191 of PK (1977) are kept, although some of them, like for the aquifer
 192 thickness, H_{1i} , in Fig. 1(a), may look bizarre to groundwater
 193 hydrologists. Eqs. (10–34) and (10–35) in PK (1962, 1977) are
 194 written for a general case of a nonempty tailwater [Fig. 1(b)].

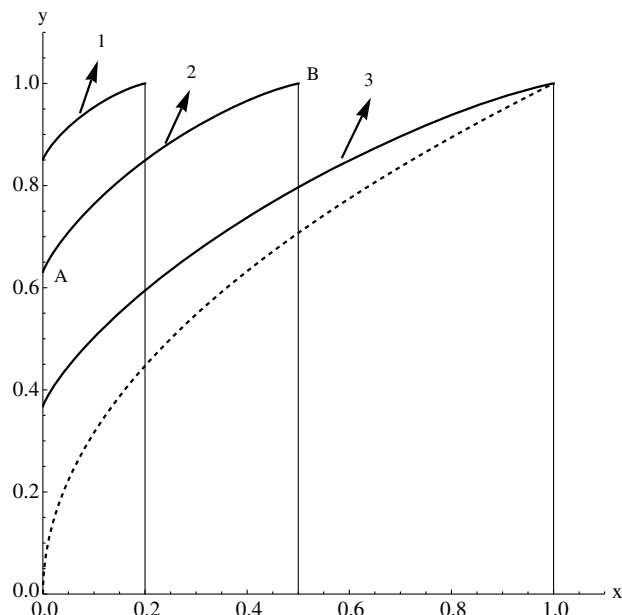
195 Correspondingly, they contain another parameter α , which is zero
 196 for this case and hence vanished in Eq. (5).

197 At the time of PK's work, determination of the two parameters
 198 (α , β) by solving a system of nonlinear equations with integrals
 199 whose integrands were special (elliptic) functions was prohibitively
 200 complicated. So, PK (1962) presented some asymptotic expansions
 201 of integrals and in PK (1977) even these expansions were dropped.
 202 Neither PK (1977, 1962) contain a systematic analysis of the shape
 203 of AB in Fig. 1(b). Hornung and Krueger (1985) extended the PK
 204 (1962, 1977) analysis and presented numerical results for several
 205 l/H_1 values in Fig. 1(b). Their motivation was "Though Polubar-
 206 inova-Kochina published her formulas in 1962, her solution was
 207 seldomly used as a reference to test numerical methods. This
 208 may be due to the fact that the evaluation of these formulas is
 209 not straightforward." The solution of the dam problem was pub-
 210 lished in the 1930s; half a century later, geotechnical engineers
 211 did not use the PK solution and spurred the Hornung and Krueger
 212 (1985) analysis; 30 years after their paper the situation is the same:
 213 with all the juggernauts of FEFLOW, HUDRUS2D, and MOD-
 214 FLOW, the PK (1962, 1977) results are not in the arsenal of numeri-
 215 cal modelers and practitioners.

216 Nowadays, solving Eq. (5) or a system of equations for (α , β),
 217 i.e., for the most general case of the dam problem with an arbitrary
 218 tailwater level in Fig. 1(b), is a routine of Wolfram's (1991) *Math-*
 219 *ematica* (or other computer algebra packages like *MATLAB*). The
 220 FindRoot, EllipticK, and NIntegrate built-in functions of *Mathema-*
 221 *tica* were used and Eq. (5) was solved as $\beta = \beta(l/H_1)$. Then
 222 Eq. (10.37) was used for determining the size of the seepage face
 223 in Fig. 1(b)

$$H_0^* = \frac{H_{0i}}{H_{1i}} = \frac{H_0}{H_1} = \frac{\int_0^{\pi/2} \frac{K[\cos^2 \chi] \sin \chi}{\sqrt{1-(1-\beta) \sin^2 \chi}} d\chi}{\int_0^{\pi/2} \frac{K[\beta+(1-\beta) \sin^2 \chi]}{\sqrt{\beta+(1-\beta) \sin^2 \chi}} d\chi} \quad (6)$$

224 The corrected PK parametric equations of BA follow from
 225 Eq. (10.41):



224 **Fig. 2.** Shapes of phreatic surface $y(x)$ in Fig. 1(b) for $l^* = 0.2, 0.5,$
 225 and 1.0 (Curves 1–3, respectively, solid lines) and the DF parabola
 (dashed line) for $l^* = 1.0$

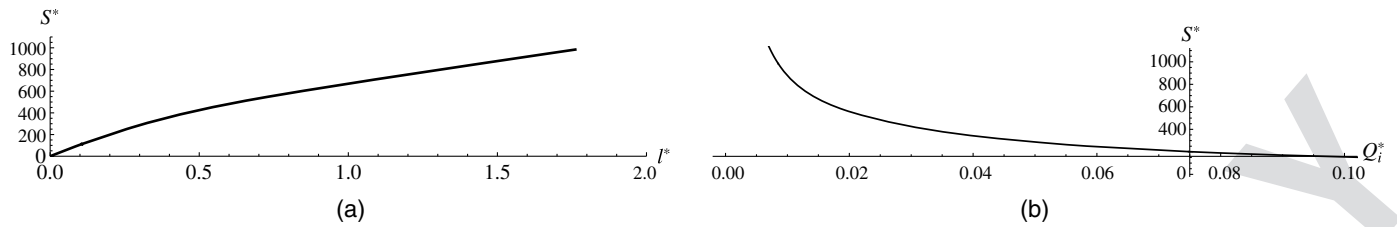


Fig. 3. (a) Dimensionless cross-sectional area S^* of the wedge as a function of the dimensionless wedge base l^* for $\delta = 0.03$; (b) wedge area as a function of the incident gradient

$$x^*(\theta) = -\frac{x_i}{H_{1i}} = \frac{x}{H_1} = \frac{l}{H_1} - \frac{\int_0^\theta \frac{K[\sin^2\chi] \sin\chi}{\sqrt{1-\beta\sin^2\chi}} d\chi}{\int_0^{\pi/2} \frac{K[\beta+(1-\beta)\sin^2\chi]}{\sqrt{\beta+(1-\beta)\sin^2\chi}} d\chi},$$

$$y^*(\theta) = -\frac{y_i}{H_{1i}} = \frac{y}{H_1} = \frac{H_0}{H_1} + \frac{\int_0^\theta \frac{K[\cos^2\chi] \sin\chi}{\sqrt{1-\beta\sin^2\chi}} d\chi}{\int_0^{\pi/2} \frac{K[\beta+(1-\beta)\sin^2\chi]}{\sqrt{\beta+(1-\beta)\sin^2\chi}} d\chi},$$

$$0 \leq \theta \leq \pi/2 \quad (7)$$

Superscripts in x and y are dropped for the sake of brevity. Then Eq. (4) is rewritten in a dimensionless form

$$S^* = \frac{S_i}{H_{1i}^2} = \frac{1}{\delta^2} \left[\frac{l}{H_1} - \int_0^{\pi/2} y(\chi) \frac{dx(\chi)}{d\chi} d\chi \right] \quad (8)$$

where $dx(\chi)/d\chi$ is evaluated from the first equation in Eq. (7).

Fig. 2 shows $y(x)$ for $l^* = l/H_1 = 0.2, 0.5,$ and 1.0 (Curves 1–3, respectively), i.e., in a benign context of the dam problem of Fig. 1(b). Table 2 of Hornung and Krueger (1985) was also checked and *Mathematica* gave exactly the same H_0/H_1 values. For comparison, at $l^* = 1.0$ a DF parabolic phreatic surface $y = \sqrt{x}$ is also plotted in Fig. 2 as a dashed line. For the selected values of l^* in Fig. 2, the DF approximation is not appropriate.

Fig. 3(b) uses the same Eqs. (8) and (3) to depict the area of the SWI zone in the context of SWI management. Fig. 3(b) shows a graph of $S^*(Q_i^*)$, where $Q_i^* = Q_i/(k_i H_{1i})$ is the uniform hydraulic gradient upstream of the SWI zone [compare with a relevant Fig. 3(a) of Ferguson and Gleeson (2012)]. At $Q_i^* \rightarrow 0$, the whole aquifer in Fig. 1(a) is occupied by seawater, i.e., the curve in Fig. 3(a) goes up to the left. The recent alarmism about the rise of the global

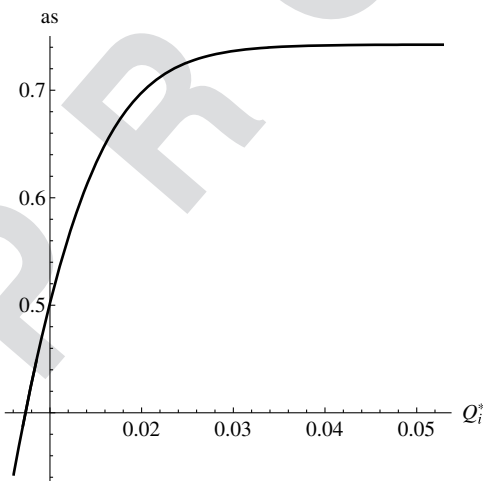


Fig. 4. Aspect ratio of the vertical to horizontal sizes of the SWI wedge in Fig. 1(a)

seawater level (H_{1i}) pedals mostly the ensued damage to on-shore structures, although Fig. 3(b) illustrates the invisible tongue extension deep inland, with a potential deleterious impact on agricultural land that is irrigated from coastal aquifers.

Fig. 4 shows $as = (H_{1i} - H_{oi})/l_i$ as a function of Q_i^* . The ratio as [the vertical size of the SWI wedge in Fig. 1(a) to its horizontal size] quantifies the degree of the hydrodynamic push of the wedge by flowing groundwater. In the case of no SWI $Q_i^* \rightarrow \infty, \beta \rightarrow 0$, the area of the wedge and both its sizes approach zero but the aspect ratio $as \rightarrow 8Ca/\pi^2 \approx 0.74$, where Ca is the Catalan constant (see the horizontal asymptote in Fig. 4), as it should be according to PK (1977) in the dam problem (see PK's Case 2).

Comparison with Sand Tank Experiments

GC conducted experiments (see their Fig. 2) in a sand-filled tank with the following values: $k_i = 1,050$ m/day; $Q_i = 1.42/2.7$ cm²/s; $H_{1i} = 26$ cm; $l_i = 15$ cm; $r_f = 1$ g/cm³; and $r_s = 1.026$ g/cm³. Although GC's experimental flow was unconfined as in Kashef (1983), i.e., instead of the caprock $E_{1i}C_iD_i$ in Fig. 1(a), GC had a phreatic surface, the slope of this surface was relatively small. In numerical modeling GC used a confined flow-transport model. The GC numerical and experimental results matched well. Therefore, the replacement of GC's free surface by a horizontal no-flow caprock as in Fig. 1(a) is reasonable for the selected experimental setup.

The theoretical value for GC's experiment, according to Eq. (3), is $Q_i = 6.15$ m²/day. GC's measured discharge is $Q_i = 4.54$ m²/day. GC's experimental value was also used for l_i in the left-hand side of Eq. (5) and the root of this equation was found to be $\beta = 0.58$. Then this β was put into the right-hand side of Eq. (6) and $H_{oi} = 10.93$ cm was found, while the GC size of the discharge window was $H_{oi} = 13$ cm.

Bertorelle (2014) conducted similar experiments and SUTRA-based numerical modeling for a sandbox with $k_i = 1.8 \times 10^{-3}$ m/s, $Q_i = 2.5/0.3 \times 10^{-3}$ m²/h, $H_{1i} = 41$ cm, $r_f = 1$ g/cm³, and $r_s = 1.025$ g/cm³. Now her experimental data are converted into dimensionless format. The theoretical formula Eq. (3) gives $Q_i^* = 0.004$, while the Italians measured the discharge of $Q_i^* = 0.0027$. The FindRoot routine was again used to solve Eq. (5) and β was found to be 0.9999334. This was used in Eq. (6), which gave a theoretical $H_0^* = 0.17$, which agrees well with Fig. 9.28 of Bertorelle's experiment and numerical modeling. Therefore, both GC's and Bertorelle's (2014) results match well the theory presented in this paper.

Conclusion

Steady SWI with a sharp interface in a confined aquifer is mathematically equivalent to the classical PK problem of a phreatic

289 surface seepage through an earth dam. This mathematical common-
290 ality is well known since the comparisons of the Glover (1959) SWI
291 problem, having a parabolic sharp interface, with the Pavlovsky
292 26 problem (PK) of flow toward a Zhukovsky drain, which has a para-
293 bolic phreatic surface.

294 With modern computer algebra tools and rectification of PK's
295 typos and errors, the analytical solution to the dam problem is meta-
296 morphosed into solution to a SWI problem, modulo stretching-
297 rescaling. The sharp-interface model matches well the experiments
298 of GC and Bertorelle (2014), as well as their numerical modelling
299 by variable density codes SEAWAT and SUTRA.

300 Fig. 3(b) corroborates the results obtained in Kacimov et al.
301 (2009) in terms of the DF model for an unconfined coastal aquifer,
302 viz, SWI increases rapidly with the decrease of the incident gra-
303 dient [uniform in the left element of Fig. 1(a)] when the gradient
304 is relatively small. A similar conclusion was drawn by Ferguson
305 and Gleeson [2012, Fig. 3(a)]. For example, from Fig. 3(b), with
306 the decrease of the incident gradient Q_i^* from 0.033 to 0.0083 the
307 dimensionless area of the nasty SWI wedge in Fig. 1(a) increases
308 from 400 to 1,000. Unfortunately, in the Gulf countries a contin-
309 uing overabstraction of fresh groundwater from deeper and deeper
310 aquifers, which submarinely discharges into the sea, results in a
311 drastic SWI. The wedge encroachment inland is pretty limited
312 when the incident fresh groundwater gradient is above a certain
313 threshold level; below it a catastrophic expansion of the SWI zone
314 takes place.

315 While phreatic coastal aquifers can be replenished by relatively
316 cheap managed aquifer recharge schemes, like infiltration basins,
317 the fate of deep confined aquifers is bleak because these aquifers
318 require more expensive well injection for recuperation of SWI.

319 The sad fact of a highly nonlinear nastiness of the wedge size,
320 evidenced in the increase of the curve in Fig. 3(b) at small incident
321 fresh groundwater gradients caused either by droughts overpump-
322 ing (decrease of recharge Q_i) or increase of seawater level H_i , has
323 to win the hearts and minds of water resource managers in the Gulf
324 and other SWI-prone regions. Hopefully groundwater engineers
325 will sympathize with the authors' predilection for analytical solu-
326 tions, in particular, the old PK one for the dam problem, which was
327 exploited in this paper.

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