

WEYL GEODESIC FIELD OF CONES
IN A THREE-DIMENSIONAL RIEMANNIAN SPACE

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The notion of a geodesic field of cones of directions in a Riemannian space was introduced in [1]. A field of cones is defined with the use of a symmetric tensor $a_{ij}(x^k)$ (x^k are coordinates in the space; $i, j, k = \overline{1, n}$; n is the dimension of the space) by the equation

$$a_{ij}X^iX^j = 0, \tag{1}$$

where X^i are the components of a contravariant vector determining a direction. A field of cones (1) is called geodesic if all geodesics whose directions coincide at some point with a direction satisfying (1) retain this property at all their points. As was shown in [1], the following equalities are a criterion for a field of cones (1) to be a geodesic one: $\nabla_{(k}a_{ij)} = M_{(k}a_{ij)}$, where M_k are the components of some vector, ∇_k is the symbol of covariant differentiation with respect to the connection under consideration, and the brackets denote the symmetrization with respect to the corresponding indices.

In [2], fields of cones (1) in a Riemannian space were considered whose nondegenerate tensor together with some vectors M_k and R_k and the metric tensor g_{ij} satisfy the equation

$$\nabla_{(k}a_{ij)} = M_{(k}a_{ij)} + R_{(k}g_{ij)}. \tag{2}$$

In [3], such fields were called Weyl geodesic. Riemannian spaces admitting Weyl geodesic fields of cones generated by a nondegenerate tensor field a_{ij} of a special type whose characteristic equation

$$|a_{ij} - \lambda g_{ij}| = 0 \tag{3}$$

has at least two real roots were constructed in [4].

In this paper, we construct three-dimensional Riemannian spaces admitting a special (see § 2 below) Weyl geodesic field of cones whose characteristic equation has a pair of complex conjugate roots or a real root of multiplicity 3 (the latter case is excluded in [4]).

1. Weyl geodesic field of cones

Consider some properties of Weyl geodesic fields of cones.

First of all, the following should be noted: although the tensor of such a field is defined up to an arbitrary functional factor, the geometrical properties of this field do not depend on the choice of a normalization of its tensor.

In fact, for a tensor \tilde{a}_{ij} such that $a_{ij} = \alpha\tilde{a}_{ij}$ (α is a function of the coordinates), from (2) it follows that $\nabla_{(k}\tilde{a}_{ij)} = \tilde{M}_{(k}\tilde{a}_{ij)} + \tilde{R}_{(k}g_{ij)}$, where $\tilde{M}_k = M_k - \alpha^{-1}\partial_k\alpha$, $\tilde{R}_k = \alpha^{-1}R_k$, $\partial_k \equiv \frac{\partial}{\partial x^k}$. In addition, the form of equations (2) is invariant with respect to conformal transformations ([5], p. 161) of connection.